# Observations of the Prompt Gamma-Ray Emission of GRB 070125

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## ABSTRACT

The long, bright gamma-ray burst GRB 070125 was localized by the Interplanetary Network. We present light curves of the prompt gamma-ray emission as observed by Konus-WIND, RHESSI, Suzaku-WAM, and Swift-BAT. We detail the results of joint spectral fits with Konus and RHESSI data. The burst shows moderate hard-to-soft evolution in its multi-peaked emission over a period of about one minute. The total burst fluence as observed by Konus is  $1.79 \times 10^{-4} \text{ erg/cm}^2$  (20 keV-10 MeV). Using the spectroscopic redshift z = 1.548, we find that the burst is consistent with the "Amati"  $E_{peak,i} - E_{iso}$  correlation.

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Assuming a jet opening angle derived from broadband modeling of the burst afterglow, GRB 070125 is a significant outlier to the "Ghirlanda"  $E_{peak,i} - E_{\gamma}$  correlation. Its collimation-corrected energy release  $E_{\gamma} = 2.5 \times 10^{52}$  ergs is the largest yet observed.

Subject headings: gamma-rays: bursts

#### 1. Introduction

The prompt gamma-ray emission of gamma-ray bursts (GRBs) is the most extensively studied aspect of these energetic explosions. Indeed, for twenty-five years after the discovery of GRBs (Klebesadel et al. 1973), the prompt emission was the only GRB observable available. With the first afterglow observations at longer wavelengths (Costa et al. 1997; van Paradijs et al. 1997), detailed analysis of burst models became possible. Presently, the Swift satellite is detecting  $\sim 100$  bursts per year, most with rapid localization and followup.

The exact mechanism which produces the prompt gamma-ray emission, with its characteristic smoothly broken power-law spectrum, has not been definitively established. Recent efforts to correlate burst observables with the intrinsic burst energetics have increased the importance of detailed spectral fitting for localized bursts (for a review, see Zhang 2007). Some correlations involve the peak spectral energy  $E_{peak}$ , which is often above the  $\sim 150 \text{ keV}$  cutoff of the Swift Burst Alert Telescope (BAT) passband.

Several current observatories are capable of detailed spectral analysis of GRBs over the full range of  $E_{peak}$ . Konus-W (Aptekar et al. 1995) is a double scintillator instrument on the WIND spacecraft. The Ramaty High Energy Solar Spectroscopic Imager (RHESSI) is a solar observatory which uses nine germanium detectors to image the Sun at X-ray to gammaray energies (Lin et al. 2002). RHESSI's detectors are unshielded and receive emission from astrophysical sources like GRBs. The Wide-Band All-Sky Monitor (WAM) (Yamaoka et al. 2005) aboard Suzaku is the large BGO anticoincidence shield for the Suzaku Hard X-Ray Detector. AGILE (Tavani et al. 2006) and GLAST (Ritz 2007) will give additional coverage at the energy range of  $E_{peak}$  and extend spectral coverage for GRBs up to tens of GeV.

In this paper, we present Konus, RHESSI, and Suzaku observations of the bright GRB 070125. In Section 2, we discuss the observations and the localization of the burst by the IPN. Section 3 contains the burst light curves, and in Section 4 we conduct joint spectral fits to the Konus and RHESSI data.

### 2. Observations

GRB 070125 was observed by six spacecraft in the Interplanetary Network (IPN): RHESSI, Suzaku WAM, and Swift-BAT, all in low Earth orbit; the anticoincidence system of the spectrometer aboard the International Gamma-Ray Astrophysics Laboratory (INTE-GRAL), at 0.44 light-seconds from Earth; Konus-Wind, at 5.4 light-seconds from Earth; and the High Energy Neutron Detector and Gamma Sensor Head aboard Mars Odyssey 2001, at 1130 light-seconds from Earth. The two other distant missions in the network, Ulysses and MESSENGER (Mercury Surface, Space Environment, Geochemistry, and Ranging), were off. Since Swift was slewing at the time of the burst, it did not immediately localize it. However, the source appeared in a routine image made after the slew was completed, and its 2.5' radius error circle was consistent with the initial IPN localization (Hurley et al. 2007). Even with more than six minutes of elapsed time since the burst onset, the BAT image detections were highly significant at 8.2 sigma (Racusin et al. 2007).

With only one distant spacecraft, the IPN localized the event to a long, narrow error ellipse whose area  $(3\sigma)$  is  $\sim 1200$  square arcminutes, centered at RA(2000) = 07h 51m 17.85s,  $Dec(2000) = +31^{\circ} 06' 12.78"$ . The chi-squared for this position is 1.57 for 3 degrees of freedom. Figure 1 shows the central region of the error ellipse, with the BAT 90% confidence error circle and the optical counterpart.

Initial spectral fits to the prompt emission were reported for RHESSI by Bellm et al. (2007) and for Konus by Golenetskii et al. (2007). The initial RHESSI best fit model was a cutoff power law (equivalent to the Band function below  $E_{break}$ , see §4) with  $\alpha=1.33^{+0.11}_{-0.09}$ ,  $E_{peak}=980.\pm300.$  keV, and a 30 keV–10 MeV fluence of  $1.5\times10^{-4}$  erg/cm<sup>2</sup>. The Konus data were best fit by a Band function with  $\alpha=-1.10^{+0.10}_{-0.09}$ ,  $\beta=-2.08^{+0.10}_{-0.15}$ ,  $E_{peak}=367^{+65}_{-51}$  keV. The measured Konus 20 keV–10 MeV fluence was  $(1.74^{+0.18}_{-0.15})\times10^{-4}$  erg/cm<sup>2</sup>. All errors are 90% C.L.

Pelangeon and Atteia derived a pseudo-redshift for this burst by using the RHESSI parameters (Pelangeon & Atteia 2007b) and the Konus values (Pelangeon & Atteia 2007a). These were fairly consistent at  $1.6 \pm 0.8$  and  $1.3 \pm 0.3$  respectively.

Cenko & Fox (2007) reported an optical counterpart at RA(2000)=07h 51m 17.75s,  $Dec(2000) = +31^{\circ}$  09' 04.2". This counterpart was confirmed by Updike et al. (2007) in the R band.

Racusin & Vetere (2007) reported detection by the *Swift* XRT. The XRT position was RA(J2000) = 7h 51m 18.08s,  $Dec(J2000) = +31^{\circ} 09' 02.2''$ , 4.7 arcseconds from the optical transient reported by Cenko & Fox (2007).

Initial afterglow detections in other bands included *Swift* UVOT in the UV (Marshall et al. 2007), radio (van der Horst 2007), and IR (Bloom et al. 2007). Milagro (Dingus 2007) observations of the source took place, but no VHE gamma-ray source was detected.

Fox et al. (2007) reported a redshift of  $z \ge 1.547$  for GRB 070125 from the identification of the Mg II doublet. Cenko et al. (2008) tightened this estimate to  $z = 1.5477 \pm 0.0001$ . Independent observations by Prochaska et al. (2007), reported by Updike et al. (2008), reveal absorption features which are consistent with z = 1.548 if identified as C IV and Si IV, and the absence of Lyman absorption features requires z to be near this value.

Observations of the decaying afterglow yielded multiple possibilities for a jet break. The Swift-XRT data showed a possible jet break at  $1.35 \pm 0.35$  days, but were also consistent with no jet break (Burrows & Racusin 2007; Racusin et al. 2007). Independent optical observations (Mirabal et al. 2007; Garnavich et al. 2007) showed a break in the decay at  $t \geq 4$  days. The non-detection by Chandra (Cenko et al. 2007) was also consistent with a break occurring after 4 days. Updike et al. (2008) used a larger optical dataset to fit a jet break time of  $t = 3.73 \pm 0.52$  days, but cautioned that flaring made the best fit break time dependent on the choice of time intervals. Chandra et al. (2008) found a best fit break time of t = 3.8 days in a joint optical–X-ray fit. They suggested that the break might be chromatic, as the X-ray data alone did not require a break, and proposed that inverse Compton emission could create a delay between the optical and X-ray breaks.

Extensive observations of the afterglow of GRB 070125 allowed detailed studies of the unusual burst environment. Cenko et al. (2008) suggested that the low absorbing column densities inferred from the afterglow spectra indicate that this long burst took place in a low-density galactic halo. Chandra et al. (2008) performed detailed broadband fitting of the afterglow, and concluded that the immediate environment of the progenitor was likely high density ( $n \sim 50 \text{ cm}^{-3}$  for a constant density profile). They also found evidence that the gamma-ray production efficiency for this burst was unusually high ( $\eta_{\gamma} \sim 0.65$ ).

#### 3. Light Curve

Figure 2 shows the Konus, RHESSI, Suzaku-WAM, and Swift-BAT light curves corrected for light travel time between the spacecraft. The Konus trigger time was  $T_{0,KW}=07:20:50.853$ . Photon travel time from RHESSI to Konus was 5.197 seconds, from Suzaku to Konus was 5.202 seconds, and from Swift to Konus was 5.215 seconds.

The light curves show a qualitatively similar multi-peaked structure with roughly four major periods of emission. The RHESSI data in interval A have a slight but significant

feature around  $T_0 + 4$  seconds whose origin is unclear. The bump appears in data from all three detectors used in this study. Examination of hardness ratios suggests that the bump is softer than the rest of the emission in the interval, but insignificantly so ( $\sim 1\sigma$ ). The difference is even more negligable when we consider only data above 65 keV. Accordingly, the bump (if extraneous) should not meaningfully influence the spectral fits reported in Section 4.

T90 for the Konus light curve was  $62.2 \pm 0.8$  seconds (20–1150 keV), for RHESSI  $63.0 \pm 1.7$  seconds (30 keV–2 MeV), and for Suzaku  $55 \pm 2$  seconds (50 keV–5 MeV). In the individual Konus bands, the T90s were  $62.8 \pm 1.8$  seconds (G1: 20 keV–75 keV),  $61.5 \pm 0.9$  seconds (G2: 75 keV–300 keV), and  $60.0 \pm 5.6$  seconds (G3: 300 keV–1150 keV). Uncertainties on all T90s are 1-sigma and were obtained by perturbing the light curves with Poisson noise and finding the new T90 values for 1000 trials. Racusin et al. (2007) report a T90 of 60 seconds for the *Swift*-BAT light curve. Because *Swift* did not trigger on the burst, no BAT event data were stored. The available rate data contain slew artifacts; accordingly, we do not perform further analysis on the BAT data.

Both Konus and RHESSI observed the 64 millisecond peak flux at  $T-T_0 = 41.472$  seconds. Using the spectral fits from Section 4, the peak flux (20 keV–10 MeV) observed by Konus was  $(1.85^{+0.35}_{-0.36}) \times 10^{-5}$  erg/cm<sup>2</sup>/s. RHESSI observed a peak flux of  $(2.92^{+0.68}_{-0.63}) \times 10^{-5}$  erg/cm<sup>2</sup>/s. While the RHESSI fluences computed in Section 4 are lower than those measured by Konus, RHESSI recorded a greater proportion of counts in the 64 ms peak interval, implying a larger peak flux. These values are moderately sensitive to background subtraction; the errors quoted are purely statistical.

Figure 3 shows the fast time evolution of hardness ratios for Konus and Suzaku. The burst shows a general softening trend in time, excepting the period of peak flux in interval C, which has comparable hardness to the initial emission in interval A.

#### 4. Spectral Analysis

We performed spectral analysis for the time intervals given in Table 1 using the Konus and RHESSI data. While spectral data are available from Suzaku, the GRB photons passed through the X-Ray Spectrometer (XRS) dewar before reaching the WAM. This direction is not well-calibrated for the WAM, in part due to uncertain levels of solid Ne in the dewar. With the detector response poorly understood, it is impossible to determine effectively the spectral parameters. Accordingly, we omit the Suzaku data in the spectral fits.

Konus 64-channel spectra are available beginning 0.512 seconds before the trigger and

are integrated over variable timescales. The detector response, which is a function only of the burst angle relative to the instrument axis, is generated from Monte Carlo simulations described by Terekhov et al. (1998).

Because of radiation damage to the RHESSI detectors, only three of the nine detectors (rear segments 1, 7, and 8) were usable for this analysis. While the damaged detectors continue to record significant counts, the effect of the radiation damage on the spectral response has proven difficult to model.

To generate the RHESSI spectral response, we simulated monoenergetic photon beams impinging on a detailed mass model in the Monte Carlo suite MGEANT (Sturner et al. 2000). The response of each detector changes as RHESSI rotates, so we used a beam geometry with photons generated along 60° arcs in rotation angle. The resulting sector responses were weighted by the burst light curve and added together. Fit results were not appreciably different when using a simple azimuthally averaged response. The beam made an angle of 165° with the RHESSI rotation axis to match the off-axis angle of the GRB (165.2 degrees). The simulated photons had initial energies given by 192 logarithmically-spaced bins from 10 keV to 30 MeV.

We conducted the spectral fitting in parallel using the spectral fitting packages XSPEC v11<sup>1</sup> and ISIS v1.4.3 (Houck 2002). The fit parameters obtained from both programs were identical. Robust fitting required a lower fit bound of 65 keV for RHESSI, slightly higher than the typical 30 keV lower limit. Because the GRB was arriving from the extreme rear of RHESSI, the photons passed through the back plate of the RHESSI cryostat and were hence subject to greater attenuation at low energies. The fit ranges were accordingly 20 keV–10 MeV for Konus and 65 keV–10 MeV for RHESSI. We rebinned the data to a minimum S/N of 2 before performing the spectral fits. This rebinning did not greatly affect the best fit parameters. Fluence errors were obtained in ISIS by stepping through a grid of fluence values, refitting the free parameters at each grid point, and monitoring the change in chisquared. Since it does not assume that the statistic space is quadratic, this method provides more accurate values for the uncertainties than those generated in XSPEC with the flux command.

The data were well-fit in intervals A-C by a Band function (Band et al. 1993):

$$N_E = \begin{cases} A(E/E_{piv})^{\alpha} \exp(-E(2+\alpha)/E_{peak}) & E < E_{break} \\ B(E/E_{piv})^{\beta} & E > E_{break} \end{cases}$$

with 
$$E_{break} \equiv E_{peak} \frac{(\alpha-\beta)}{(2+\alpha)}$$
 and  $B \equiv A(\frac{(\alpha-\beta)E_{peak}}{(2+\alpha)E_{piv}})^{\alpha-\beta} \exp(\beta-\alpha)$ . For  $\beta < -2$  and  $\alpha >$ 

<sup>&</sup>lt;sup>1</sup>http://heasarc.gsfc.nasa.gov/docs/xanadu/xspec/

-2,  $E_{peak}$  corresponds to the peak of the  $\nu F_{\nu}$  spectrum. The normalization A has units photons/(cm<sup>2</sup> s keV), and  $E_{piv}$  is here taken to be 100 keV. For joint fits, the Band function parameters  $\alpha$ ,  $\beta$ , and  $E_{peak}$  were tied for both instruments, but the normalizations were allowed to vary independently. For interval D, the best fit model after grouping was a simple power law. We report the best-fit spectral parameters in Table 2. Figure 4 shows the spectra in all intervals for the joint fit.

For single-instrument fits, the Konus data provide superior fit quality and better constraint on the fit parameters, due in part to having about six times more usable counts. The fit fluence,  $\alpha$ , and  $\beta$  are generally consistent between RHESSI and Konus. However, the RHESSI data prefer higher  $E_{peak}$ , matching the best fit Konus values only at the lowest end of rather large error bars. The Konus fit parameters for the total burst match well the initial values reported via the GCN (Golenetskii et al. 2007). The RHESSI fit  $E_{peak}$  typically is lower here than in the value reported in the GCN (Bellm et al. 2007), but this difference is expected from fitting using the Band function rather than a cutoff power law (Band et al. 1993).

The spectral parameters for the joint fits are consistent with the Konus-only values. There are slight improvements in the uncertainties of some of the fit parameters at a cost of an increase in the chi squared. The RHESSI residuals in the joint fit (Figure 4) show a characteristic deviation pattern, indicating that the instruments disagree on the spectral shape. The Konus data dominate the fit because of their better statistical quality. The residuals for the RHESSI-only fits do not show any systematic deviation.

For intervals A, and B, the ratio of the RHESSI normalization to the Konus normalization is 0.88. For interval C, the ratio is 0.95. Characteristic uncertainties for the ratio are 0.04-0.05. In interval D, the ratio for the power-law fit is  $0.84^{+0.14}_{-0.13}$ . Absolute normalizations in photons/(cm<sup>2</sup> s keV) using  $E_{piv}=100$  keV for the total interval were  $(2.50^{+0.18}_{-0.15})\times 10^{-2}$  (Konus) and  $(2.25^{+0.16}_{-0.14})\times 10^{-2}$  (RHESSI).

The time-resolved fits show a moderate hard-to-soft evolution.  $E_{peak}$  is largest in the initial broad pulse (539 keV) and then softens to 355 keV in interval B. The sharp pulse in interval C has a harder spectrum (418 keV). While the statistically preferred model for the S/N grouped data in interval D is a simple power law, fitting a cutoff power-law with the Konus data to 2 MeV gives an estimate of  $E_{peak}$  at 220 keV. The high-energy spectral index  $\beta$  softens monotonically through intervals A-C.

## 5. Energetics

Knowledge of the burst redshift z = 1.548 makes it possible to draw conclusions about the overall burst energetics. We assume a standard flat cold dark matter cosmology ( $\Lambda$ CDM), with parameters ( $\Omega_{\Lambda}$ ,  $\Omega_{M}$ ,  $H_{0}$ ) = (0.761, 0.239, 73 km s<sup>-1</sup> Mpc<sup>-1</sup>), consistent with results from WMAP year 3 (Spergel et al. 2007) and large scale structure traced by luminous red galaxies (Tegmark et al. 2006). This particular set of values corresponds to the "Vanilla model" of Tegmark et al. (2006).

Extrapolating to a GRB rest-frame energy band of 1 keV-10 MeV, the isotropic emitted energy for the total burst is  $(9.59 \pm 0.39) \times 10^{53}$  ergs (Konus) and  $(8.67 \pm 0.38) \times 10^{53}$  ergs (RHESSI) for the joint fit. Because we allow independent normalizations for the Konus and RHESSI data, we obtain two values of  $E_{iso}$  from the joint fit, one for each instrument. 90% C.L. errors are obtained by exploration of the parameter space as for the fluence; we neglect uncertainty in z. These values, together with the spectral fit of the time-integrated spectrum, are consistent with the "Amati relation" correlating  $E_{iso}$  with the intrinsic peak energy of the spectrum in the GRB rest frame  $E_{peak,i}$  (Amati et al. 2002; Amati 2006; Ghirlanda et al. 2008). We plot GRB 070125 in the  $E_{peak,i} - E_{iso}$  plane in Figure 5.

Because the best fit Band function has a hard tail ( $\beta \sim -2$ ), the fluence integral is sensitive to the choice of upper energy bound. If we use the observed energy band 20 keV–10 MeV, corresponding to a GRB frame band of 50 keV–25.5 MeV, the fluence is 14% larger than that in the usual bolometric band. For consistency with previous works, we will use the 1 keV–10 MeV band for bolometric estimates.

Converting the 64 ms peak fluxes reported in Section 3 to bolometric peak luminosities using the best fit Band parameters, we find peak luminosities of  $(2.59^{+0.36}_{-0.37}) \times 10^{53}$  ergs/s for Konus and  $(4.25^{+0.87}_{-0.79}) \times 10^{53}$  ergs/s for RHESSI.

Chandra et al. (2008) performed a broadband fit to afterglow data for GRB 070125. They determined a jet opening angle of  $13.2 \pm 0.6$  degrees in their most plausible scenario (a radiative fireball expanding into a constant density (ISM) medium and emitting via synchrotron and inverse Compton channels). This jet angle was consistent with that inferred from the jet break time  $\sim 3.7$  days and an emission radius derived from radio scintillation. For the collimation-corrected energy  $E_{\gamma} = (1 - \cos \theta) E_{iso}$ , we find  $E_{\gamma} = (2.52 \pm 0.24, 2.27 \pm 0.22) \times 10^{52}$  ergs for (Konus, RHESSI). These values are the largest yet recorded for a burst with measured  $E_{peak}$  (c.f. Frail et al. 2006; Kocevski & Butler 2008 also reported lower limits on  $E_{\gamma}$  greater than  $10^{52}$  ergs for several Swift bursts using the time of the last XRT observation). We plot GRB 070125 in the  $E_{peak,i} - E_{\gamma}$  plane in Figure 6 to examine its consistency with the "Ghirlanda"  $E_{peak,i} - E_{\gamma}$  correlation (Ghirlanda et al.

2004, 2007).

In Figures 5 and 6, we also overplot the best fit correlation lines. A number of fitting approaches have been considered in the literature in an effort to account for the apparent extra-statistical spread of the points about the correlation (for a review, see Ghirlanda et al. 2008). We have followed Ghirlanda et al. (2008) in presenting two least squares fits, one in which the data points are unweighted and a second in which the errors on both axes are considered. After the fit, we estimate the dispersion of the points perpendicular to the best fit correlation line using the square root of the bias-corrected sample variance.

GRB 070125 is quite consistent with the Amati relation: including it in the fit makes negligible changes in the best-fit correlation slope or the logarithmic dispersion (0.20 dex). However, it is a  $5.0\sigma$  outlier to the Ghirlanda correlation fitted without it, using the sample dispersion to estimate  $\sigma$ . Including GRB 070125 in an unweighted fit of the bursts in the Ghirlanda sample, the overall dispersion increases to 0.13 dex (from 0.09 dex), and GRB 070125 remains a  $2.8\sigma$  outlier.

The unusual environment of GRB 070125 is responsible for its high value of  $E_{\gamma}$ . In particular, the jet opening angle of  $13.2 \pm 0.6$  degrees derived by Chandra et al. (2008) is larger than all of those presented by Ghirlanda et al. (2007). Retaining the  $3.7 \pm 0.5$  day jet break time well-established in the optical (Updike et al. 2008; Chandra et al. 2008), we may derive the jet opening angle assuming adiabatic emission and more conventional parameters (Sari et al. 1999). Assuming an ISM profile with circumburst density n=3 cm<sup>-3</sup> and a gamma-ray production efficiency of  $\eta_{\gamma}=0.2$ , the corresponding jet opening angle is  $\theta=5.6\pm0.3$  degrees for Konus. The resulting collimation-corrected energy would be  $E_{\gamma}=(4.6\pm0.5)\times10^{51}$  ergs, only  $0.8\sigma$  from the best-fit correlation omitting GRB 070125.

## 6. Discussion

While GRB 070125 had a large measured prompt gamma-ray fluence, its spectral properties are unremarkable. The values of the best-fit spectral parameters are similar to those observed for other bright bursts (e.g., Kaneko et al. 2006), and the spectral evolution observed is similarly common. The environment of GRB 070125 is unique, however (Cenko et al. 2008; Chandra et al. 2008; Updike et al. 2008), requiring a broad jet opening angle in broadband afterglow models (Chandra et al. 2008). After collimation correction, GRB 070125 has the most energetic prompt emission yet observed and is a significant outlier to the correlation between peak energy and  $E_{\gamma}$ .

GRB 070125 appears to weaken the claim that the Ghirlanda correlation has low dis-

persion. GRB 070125 is not a "recognizable" outlier to the Ghirlanda relation in the sense of Ghirlanda et al. (2007), as it is highly consistent with the Amati relation. Its jet parameters have been derived from a rich and well-sampled afterglow dataset. While the circumburst environment of this GRB is unusually dense, this only highlights the assumption of a fairly narrow range of efficiency and density parameters for the majority of GRBs where broadband modeling of the afterglow has not been possible. The true dispersion of the correlation may in fact be larger.

The physical significance of GRB spectrum–energy correlations has been questioned (e.g. Butler et al. 2007, 2008). In particular, detector trigger thresholds affect burst detection, and more complex selection effects govern the measurement of peak energies, redshifts, and afterglow breaks. These effects can influence the sample of GRBs with known redshift,  $E_{peak,i}$ , and  $E_{\gamma}$ . Ghirlanda et al. (2008) examined the effect of trigger and spectral analysis thresholds in the  $E_{peak}$ –fluence plane, finding that the Swift-detected burst sample was truncated by the spectral analysis threshold. Neither threshold truncated the pre-Swift burst sample.

We were unable to confirm the source of the systematic shift in  $E_{peak}$  and fluence between the two instruments for this burst. Minor radiation damage was becoming noticeable in RHESSI detector 8 near the time of this work, mostly below the 65 keV cut utilized here. It is also possible that the Monte Carlo simulation of the RHESSI response is less accurate for such extreme off-axis angles, where a greater number of interactions with the cryostat may be expected.

Our previous work had found excellent agreement in all fit parameters for independent RHESSI and Konus spectral fits for GRB 051103 and GRB 050717. For the short GRB 051103, Konus found  $E_{peak} = 1920 \pm 400$  keV and a 20 keV-10 MeV fluence of 4.4  $\pm$  0.5  $\times$  10<sup>-5</sup> ergs/cm<sup>2</sup> (Golenetskii et al. 2005; Frederiks et al. 2007). A RHESSI fit yielded  $E_{peak} = 1930 \pm 340$  keV and 20 keV-10 MeV fluence of 4.5  $\times$  10<sup>-5</sup> erg/cm<sup>2</sup> (Bellm et al. 2006). Krimm et al. (2006) found for a cutoff power-law fit to Konus data for GRB 050717 a best fit value of  $E_{peak} = 2101^{+1934}_{-830}$  keV. A RHESSI fit to the same burst found  $E_{peak} = 1550^{+510}_{-370}$  keV (Wigger et al. 2006). Those bursts had RHESSI off-axis angles of 97 and 110 degrees, respectively.

Joint spectral fits to *Swift*-BAT and RHESSI data for 25 bursts co-observed by the two instruments between December 2004 and December 2006 indicated that no offset in response normalization was needed for the two instruments (Bellm et al. 2008). However, for two of three bursts occurring during or after December 2006, the RHESSI data showed a significant deficit relative to *Swift*-BAT. The RHESSI polar angles for all three late bursts were between 90 and 110 degrees. These fits were conducted using only detectors 1 and 7, which do not

appear to have radiation damage in background spectra during this interval. Nonetheless, these results suggest that the observed offset in the RHESSI and Konus fit parameters found here is more likely a consequence of increased radiation damage in the RHESSI detectors than a geometric effect or a generic offset in the RHESSI simulations.

Future analysis of archival bursts may help identify the source of any systematic effects present here. It is clear, however, that joint fits between instruments capable of constraining the full range of  $E_{peak}$  are valuable in providing the most accurate and precise determination of the fit parameters.

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Facilities: RHESSI, WIND (Konus)

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Table 1. Intervals used for spectral fitting in §4. The reference time for Konus is  $T_{0,KW}=07:20:50.853$ . For RHESSI,  $T_{0,R}=T_{0,KW}$  - 5.197 s.

Interval	$T_i$ - $T_0$ (s)	$T_f - T_0$ (s)	$T_f - T_i$ (s)
A	0	13.824	13.824
A1	0	22.016	22.016
В	22.016	34.560	12.544
$\mathbf{C}$	34.560	50.432	15.872
D	50.432	66.816	16.384
AD	0	66.816	66.816

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Table 2. Best fit parameters for the Band function for Konus (K) and RHESSI (R). Errors are quoted at the 90% confidence level. For joint fits (KR), the Konus fluence is listed first. For interval D, the fit and quoted fluence are for a power-law model.

Instruments	α	β	$E_{peak} $ (keV)	20 keV-10 MeV Fluence $(10^{-5} \text{ erg/cm}^2)$	$\chi^2/\mathrm{dof}$		
Total Burst (Intervals A1-D)							
K R KR	$\begin{array}{c} -1.09^{+0.09}_{-0.08} \\ -0.90^{+0.46}_{-0.28} \\ -1.13^{+0.09}_{-0.08} \end{array}$	$\begin{array}{c} -2.09^{+0.10}_{-0.15} \\ -2.24^{+0.20}_{-0.44} \\ -2.08^{+0.09}_{-0.14} \end{array}$	$373_{-51}^{+66}  533_{-171}^{+261}  430_{-61}^{+80}$	$17.2_{-1.5}^{+1.5}$ $16.2_{-1.7}^{+1.7}$ $17.9_{-1.3}^{+1.3}$ $16.1_{-1.1}^{+1.1}$	63/60 = 1.05 $38/30 = 1.30$ $130/93 = 1.40$		
Interval A							
K R KR	$\begin{array}{c} -0.89^{+0.18}_{-0.15} \\ -0.52^{+0.68}_{-0.42} \\ -0.96^{+0.14}_{-0.11} \end{array}$	$\begin{array}{c} -1.99^{+0.12}_{-0.23} \\ -2.12^{+0.18}_{-0.30} \\ -2.04^{+0.12}_{-0.17} \end{array}$	$447_{-99}^{+154} 512_{-157}^{+249} 539_{-113}^{+129}$	$6.20_{-0.75}^{+0.71}  5.49_{-0.71}^{+0.69}  6.22_{-0.55}^{+0.56}  5.48_{-0.47}^{+0.47}$	65/61 = 1.08 $45/32 = 1.43$ $123/96 = 1.28$		
Interval B							
K R KR	$\begin{array}{c} -1.07^{+0.14}_{-0.11} \\ -1.12^{+0.60}_{-0.28} \\ -1.11^{+0.12}_{-0.11} \end{array}$	$\begin{array}{c} -2.21^{+0.14}_{-0.23} \\ -2.33^{+0.30}_{-1.20} \\ -2.17^{+0.12}_{-0.20} \end{array}$	$318_{-55}^{+65} \\ 556_{-236}^{+362} \\ 355_{-56}^{+78}$	$4.88_{-0.53}^{+0.54}  4.82_{-0.71}^{+0.68}  5.18_{-0.47}^{+0.45}  4.61_{-0.41}^{+0.40}$	45/54 = 0.84 $18/30 = 0.60$ $76/87 = 0.88$		

Table 2—Continued

Instruments	α	β	$E_{peak}$ (keV)	20 keV-10 MeV Fluence $(10^{-5} \text{ erg/cm}^2)$	$\chi^2/{ m dof}$	
Interval C						
K R KR	$-0.92^{+0.80}_{-0.34}$	$\begin{array}{c} -2.19^{+0.19}_{-0.42} \\ -2.61^{+0.54}_{0.54} \\ -2.19^{+0.17}_{-0.35} \end{array}$		$4.09_{-0.66}^{+0.64} \\ 4.02_{-0.68}^{+0.78} \\ 4.28_{-0.56}^{+0.54} \\ 4.05_{-0.51}^{+0.50}$	44/54 = 0.82 30/28 = 1.08 87/85 = 1.03	
Interval D						
K R KR	$\begin{array}{c} -1.90^{+0.07}_{-0.08} \\ -1.97^{+0.20}_{-0.25} \\ -1.87^{+0.08}_{-0.08} \end{array}$	- - -	- - -	$1.73_{-0.24}^{+0.28}$ $1.34_{-0.29}^{+0.40}$ $1.84_{-0.29}^{+0.36}$ $1.48_{-0.24}^{+0.26}$	32/42 = 0.77 15/15 = 1.03 46/57 = 0.81	

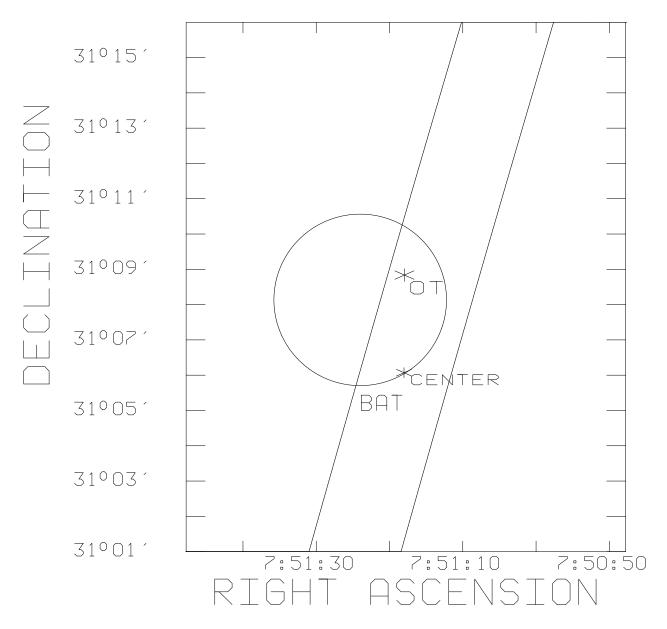


Fig. 1.— The central region of the IPN error ellipse, showing the 2.5' radius BAT  $3\sigma$  error circle, the optical transient source first reported by Cenko & Fox (2007), and the center of the ellipse. The optical source lies 0.048 degrees from the center of the IPN ellipse, on the 87% confidence contour.

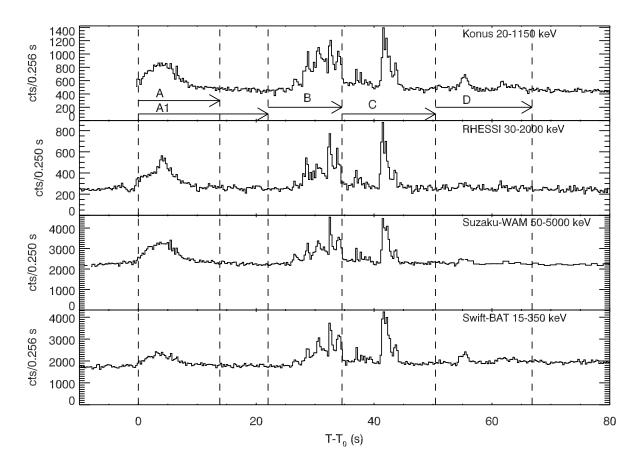


Fig. 2.— GRB 070125 light curve for Konus, RHESSI (rear segments 1, 7, and 8 only), Suzaku-WAM, and Swift-BAT. The light curves are adjusted for time of flight, with  $T_0$  given in §3. The dashed vertical lines delimit the intervals used in the time-resolved spectral fits (§4). The Swift light curve plotted contains all counts observed by Swift; in particular, it is not mask-tagged and therefore contains slew artifacts.

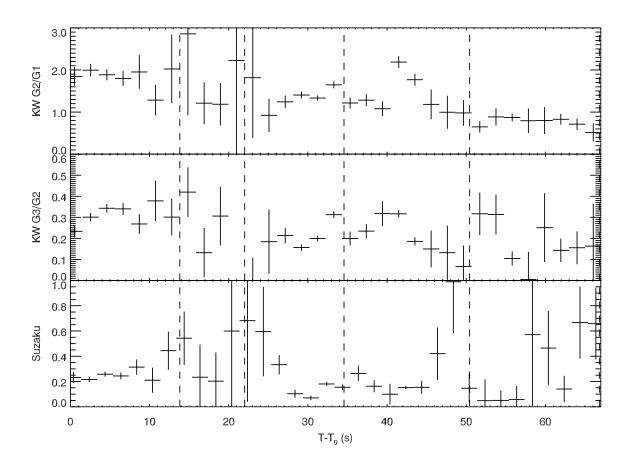


Fig. 3.— Hardness ratios for GRB 070125. The Konus energy bands are G1 (20–75 keV), G2 (75–300 keV), and G3 (300–1150 keV). The Suzaku hardness ratio plotted here is (520–5000 keV)/(50–240 keV). Dashed lines indicate spectral fitting intervals, as in Figure 2. Points near  $\sim$ 20 seconds which are off-scale for the Konus G3/G2 ratio are consistent with zero—there is negligible emission in the G3 band at this time.

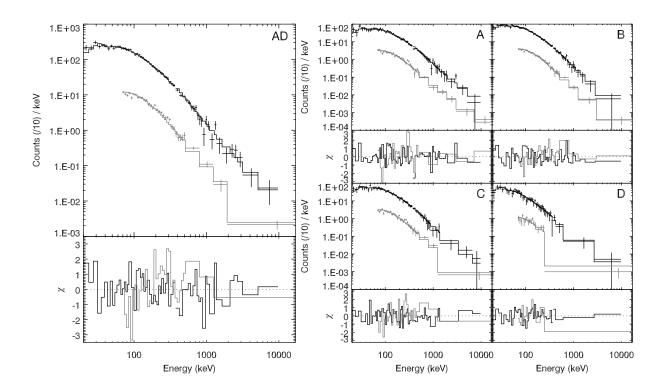


Fig. 4.— Count spectra and residuals for the joint fits. The Konus data and models are colored black, while the RHESSI data and models are gray. RHESSI data, model, and errors are divided by 10 in the count spectra plots for clarity. The overplot models differ only in normalization.

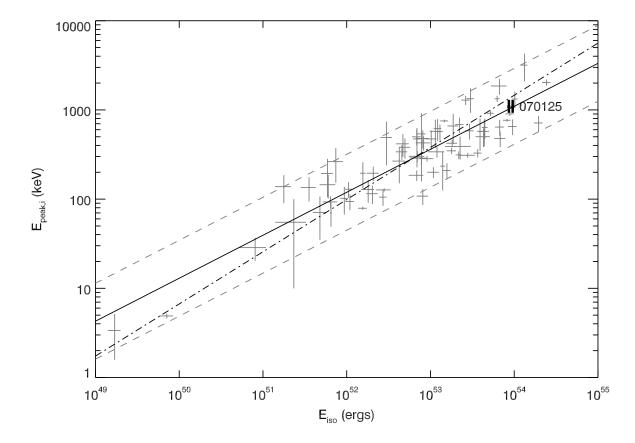


Fig. 5.—  $E_{peak,i} - E_{iso}$  correlation including GRB 070125. Values of  $E_{peak,i}$  (the intrinsic peak energy in the burst rest frame) and  $E_{iso}$  are for the joint Konus-RHESSI fit. Since the normalization was allowed to vary between the two instruments, we plot separate points for Konus and RHESSI to indicate the corresponding values of  $E_{iso}$ . The Konus data point has the larger value of  $E_{iso}$ . Data for other bursts are from Table 1 of Ghirlanda et al. (2008), plotted using the cosmology of this paper ( $\Omega_m = 0.239$ ,  $\Omega_{\Lambda} = 0.761$ , h = 0.730). The best-fit line for the unweighted data points, omitting GRB 070125, is overplot with a solid line; the  $2\sigma$  scatter about that fit is indicated with dashed lines. The dash-dotted line is the best fit when the data points are weighted by their errors on both axes, again omitting GRB 070125—see text for details.

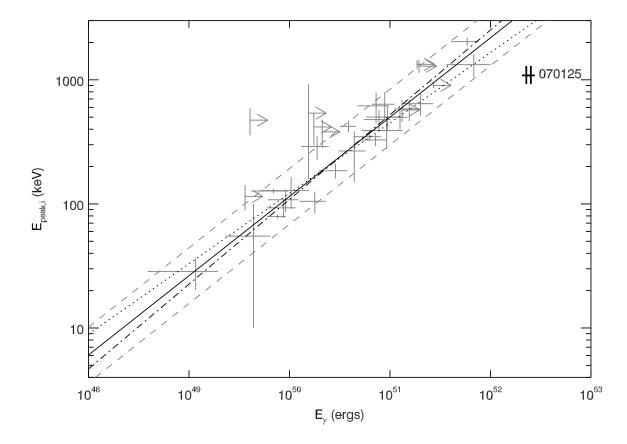


Fig. 6.—  $E_{peak,i} - E_{\gamma}$  correlation including GRB 070125. Symbols and overplot fit lines are as in Figure 5; the Konus data point has the larger value of  $E_{\gamma}$ . We also plot the best unweighted fit line including GRB 070125 with a short dotted line. Data for other bursts are from Table 1 of Ghirlanda et al. (2007), assuming an ISM density profile and plotted using the cosmology of this paper. Bursts with only lower limits on  $E_{\gamma}$  were omitted from the fit.