Abstract

Stable localised roll structures have been observed in many physical problems and model equations, notably in the 1D Swift–Hohenberg equation. Reflection-symmetric localised rolls are often found to lie on two “snaking” solution branches, so that the spatial width of the localised rolls increases when moving along each branch. Recent numerical results by Burke and Knobloch indicate that the two branches are connected by infinitely many “ladder” branches of asymmetric localised rolls. In this paper, these phenomena are investigated analytically. It is shown that both snaking of symmetric pulses and the ladder structure of asymmetric states can be predicted completely from the bifurcation structure of fronts that connect the trivial state to rolls. It is also shown that isolas of asymmetric states may exist, and it is argued that the results presented here apply to 2D stationary states that are localised in one spatial direction.
1 Introduction

To motivate the results presented in this paper, we begin with a concrete example, namely the Swift–Hohenberg equation, which is given by

$$U_t = -(1 + \Delta)^2 U - \mu U + \nu U^2 - U^3$$  \hspace{1cm} (1.1)

with $x \in \mathbb{R}$. We shall focus exclusively on the case $\mu > 0$, for which the trivial state $U = 0$ is stable, and also keep $\nu > 0$ fixed. For appropriate values of the parameters, the Swift–Hohenberg equation admits stable stationary spatially-periodic patterns which we refer to as rolls. In fact, rolls form a one-parameter family, parametrized by their wavelength. We are interested in localised roll patches, such as those shown in Figure 1, which can be thought of as subjecting rolls to a localised amplitude modulation.

The bifurcation diagrams shown in Figure 1, which have been discussed, for instance, in [3, 6, 12–14, 21, 24, 26, 29–31, 39], exhibit several intriguing features, which are commonly referred to as snaking. There are two intertwined wiggly solution branches that correspond to localised patterns whose width, measured by the $L^2$-norm, increases as we move along each branch. In particular, there exists a parameter interval so that (1.1) has infinitely many localised patterns for each $\mu$ in this interval. The end points of the interval correspond to saddle-node bifurcations of the localised states, which are aligned along a well-defined vertical asymptote. The pulses associated with the snaking branches are symmetric with respect to the reflection $x \mapsto -x$. Recently, [4, 5] found numerically that the two snaking branches discussed above are connected by a ladder branch associated with asymmetric pulses as shown in Figure 1. The asymmetric pulses bifurcate from symmetric pulses via pitchfork bifurcations.

Snaking has also been observed in the cubic-quintic Swift–Hohenberg equation

$$U_t = -(1 + \partial_x^2)^2 U - \mu U + \nu U^3 - U^5,$$  \hspace{1cm} (1.2)

see, for instance, [4, 5, 13, 14, 21, 29–31]. Equation (1.2) is equivariant under the $\mathbb{Z}_2$-action $u \mapsto -u$, and we can expect to find both even and odd localised states. This is indeed the case, and a summary of the resulting snaking observed in the references mentioned above is reproduced in Figure 1.

Planar steady states that are localised modulations in one or two spatial directions of an underlying domain-covering structure such as hexagons or rolls can also exhibit snaking, and we refer to [1, 20] for recent

Figure 1: Bifurcations diagrams of localised rolls are plotted in panels (i) and (ii) for the Swift–Hohenberg equations (1.1) with $\nu = 1.6$ and (1.2) with $\nu = 2$, respectively. The insets show localised rolls $U(x)$ as functions of $x$. Solutions on the vertical snaking branches are symmetric under the reflection $x \mapsto -x$, and the $\mathbb{Z}_2$-action $u \mapsto -u$ for (1.2), while solutions on the horizontal ladder branches are asymmetric. Ladders exist also for (1.2), see [5], but are not shown here.
Figure 2: Panel (i) shows a contour plot of a stationary solution of the planar Swift–Hohenberg equation (1.1) for \((\mu, \nu) = (0.31, 1.6)\). The solution is localised in the horizontal direction and periodic in the vertical direction. Panel (ii) shows the corresponding snaking curve. Moving up on the snaking curve causes the hexagonal region of the solution in panel (i) to expand in the horizontal direction. Reproduced from [20].

Figure 3: Localised rolls can be found as intersections of the stable and unstable manifolds of \(u = 0\) near rolls \(\gamma\) with \(\gamma(\varphi) \in \text{Fix} \mathcal{R}\). These intersections appear and disappear in fold bifurcations. As indicated here schematically, heteroclinic connections between \(u = 0\) and rolls \(\gamma\) organise the overall dynamics, and their fold bifurcations determine the asymptotic limits of the folds associated with localised rolls.

computations for the Swift–Hohenberg equation (1.1) with \(x \in \mathbb{R}^2\); see also Figure 2. Slanted snaking, where the snaking curve is not vertical but slanted to one side, has also been observed recently in [7, 10].

On a geometric level, snaking of symmetric localised 1D roll structures is well understood, and we recall briefly the key arguments from [6, 26, 39]. First, we write the steady-state equation

\[-(1 + \partial_x^2)^2 U - \mu U + \nu U^2 - U^3 = 0\]

associated with (1.1) on \(\mathbb{R}\) as the first-order system

\[u_x = f(u, \mu), \quad u = (U, U_x, U_{xx}, U_{xxx}), \quad (1.3)\]

where we chose to omit the dependence on \(\nu\), which we keep fixed. This equation has the first integral

\[\mathcal{H}(u, \mu) = u_2 u_4 - \frac{u_3^2}{2} + u_2^2 + \frac{(1 + \mu) u_1^2}{2} - \frac{\nu u_1^3}{3} + \frac{u_1^4}{4}\]

and is reversible with reverser \(\mathcal{R}u = (u_1, -u_2, u_3, -u_4)\) so that \(u(x)\) satisfies (1.3) if and only if \(\mathcal{R}u(-x)\) does. Rolls correspond to periodic orbits of (1.3), and symmetric pulses correspond to orbits \(u(x)\) homoclinic to the origin that pass close to the periodic orbits, with symmetry encoded in the condition \(u(0) \in \text{Fix} \mathcal{R}\).

The first integral allows us to restrict (1.3) to the three-dimensional zero energy level set \(\mathcal{H}^{-1}(0)\). Choosing an appropriate Poincaré section, we can further reduce the dynamics of (1.3) near a periodic orbit \(\gamma(x)\) with
zero energy to a two-dimensional map. A schematic picture of the expected dynamics of this map is shown in Figure 3. The assumption underlying Figure 3 is that heteroclinic orbits that connect the origin to the periodic orbit appear and disappear in fold bifurcations as the parameter $\mu$ is varied. Once these heteroclinic orbits are present, homoclinic orbits to the origin that pass near the periodic orbit appear also. If we plot their $L^2$-norm against the parameter $\mu$, we obtain the characteristic snaking curve shown in the centre panel of Figure 3, and we refer to [6, 39] and [20, §2.1] for geometric arguments that show why the bifurcation branch has this shape. Figure 3 shows that the folds of each homoclinic snaking curve line up with the folds of the heteroclinic orbits. The reduction outlined above can be carried out near each intersection point of $\gamma(x)$ with $\text{Fix } R$, and, since periodic orbits intersect $\text{Fix } R$ precisely twice [38], we obtain two distinct snaking curves of symmetric homoclinic orbits, in agreement with Figure 1.

The goal of this paper is to formulate a set of hypotheses about the connecting orbits between the origin and rolls that guarantee snaking. Our approach captures not only the bifurcation curves of symmetric localised structures but also of asymmetric solutions. In the remainder of the introduction, we outline the key assumptions we shall make and the conclusions we can draw from them; we refer to §2 for precise statements.

As indicated in Figure 3, we expect that the overall dynamics of snaking is organised by heteroclinic connections between the trivial state and rolls. We consider a general reversible conservative system of the form (1.3) and assume that it has a periodic orbit $\gamma(x)$ of period $2\pi$, say, with zero energy$^1$ for each $\mu$. We are interested in describing heteroclinic orbits that connect the trivial equilibrium to the periodic orbit $\gamma$. To do so, we choose a section $\Sigma_{\text{out}}$ as indicated in Figure 4 and consider intersections of the stable manifold $W^s(0, \mu)$ of $u = 0$ with the strong unstable fibre $W^{uu}(\gamma(\varphi), \mu)$ of the point $\gamma(\varphi)$ on the periodic orbit $\gamma$ in $\Sigma_{\text{out}}$. Such intersections are encoded in the set

$$\Gamma := \{(\varphi, \mu) \in S^1 \times \mathbb{R}; W^s(0, \mu) \cap W^{uu}(\gamma(\varphi), \mu) \cap \Sigma_{\text{out}} \neq \emptyset\},$$

$S^1 = [0, 2\pi]/\sim$.

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$^1$The orbit, and its period, may depend on $\mu$ but we ignore this in our outline.
and Figure 4 illustrates the additional assumption that the stable manifold $W^s(0, \mu)$ in the section $\Sigma_{\text{out}}$ is the graph of a smooth function $g(\varphi, \mu)$, locally near each heteroclinic orbit; we refer to [11] for the numerical verification of this scenario for a complex Ginzburg–Landau equation. For the sake of clarity, we focus on the case where $\Gamma$ is the graph of a function $z$,

$$\Gamma = \{(\varphi, \mu) \in S^1 \times \mathbb{R}; \mu = z(\varphi)\},$$

and examples are shown in Figure 5(i)-(ii). The set $\Gamma$ drawn in Figure 5(iii) is not a graph, and we shall outline in §6.1 how our results can be adapted to cover this situation.

Our main results assert that the complete bifurcation diagram of symmetric and asymmetric pulses can be drawn using only the function $z$, and we plot in Figure 6 the diagrams corresponding to the functions given in Figure 5. We record that a sequence of isolas of pulses, similar to the one shown in Figure 6(iii), has recently been found in [3, Figure 24] for (1.1), which suggests that fronts are present in this situation whose bifurcation structure resembles Figure 5(iii).

We shall now outline how the function $z$ determines the bifurcation diagram. We may assume that $\gamma(\varphi_0) \in \text{Fix} \mathcal{R}$ if and only if $\varphi_0 \in \{0, \pi\}$. Symmetric pulses that spend time $2L$ near the periodic orbit $\gamma$ exist if and only if

$$\mu = \mu_*(L, \varphi_0) = z(L + \varphi_0) + O(e^{-nL}), \quad L \gg 1 \quad (1.4)$$
for some $\eta > 0$, where $\varphi_0 \in \{0, \pi\}$. Furthermore, the bifurcation curves of asymmetric pulses that spend time $2L$ in the vicinity of the periodic orbit $\gamma$ are exponentially close in $L$ to the set

$\{(L, \mu) = (L, z(L + \varphi)) : z(L + \varphi) = z(L - \varphi) \text{ for some } \varphi \in S^1\}$.

This set is determined by the solutions $(L, \varphi) \in \mathbb{R}^+ \times S^1$ of

$$z(L + \varphi) = z(L - \varphi), \quad (1.5)$$

which can be found graphically as outlined in Figure 7. Using Figure 7, it is not difficult to obtain the diagrams given in Figure 6 for the sets $\Gamma$ given in Figure 5. We will also prove that the asymmetric branches that emerge from maxima of $z$ will terminate at minima. All asymmetric branches that bifurcate from the symmetric branch (1.4) with $\varphi_0 = 0$ terminate at the same branch except for the branch emerging from the global maximum of $z$, which will connect to the symmetric branch (1.4) with $\varphi_0 = \pi$, and vice versa. The start and end points of the asymmetric branches that emerge from and terminate at symmetric branches correspond to pitchfork bifurcations. Interestingly, isolas of asymmetric solutions are also possible provided the function $z$ has more than two maxima, and an example is shown in Figure 8.

Homoclinic orbits that bifurcate in reversible systems from heteroclinic cycles between equilibria have been investigated earlier in [17]. Of relevance for the snaking behaviour discussed here are homoclinic orbits that bifurcate from heteroclinic cycles connecting equilibria to periodic orbits: their geometry was elucidated in [6, 26, 39] and detailed numerical studies were carried out, for instance, in [3–5, 39]. Recently, the existence of fronts and of symmetric and asymmetric pulses was studied in [18] for the Swift–Hohenberg equation (1.1) near onset to rolls, that is near the codimension-two point $(\mu, \nu) = (0, \sqrt{27/38})$, using formal asymptotics beyond all orders. We complement these results by proving how fronts between rolls and the trivial state can be glued together to produce localised states. Our analysis is not limited to the parameter region near onset, but, in contrast to [18], it is not capable of establishing the existence of fronts in the first place. Analytical gluing results for heteroclinic cycles involving periodic orbits have, to our knowledge, been given only in the context of generic systems without reversibility. In [27], for instance, heteroclinic cycles of codimension two and three were studied, and the role played by the winding number of the stable manifold $W^s(0)$ with respect to the periodic orbit for obtaining global solution branches was illuminated. The proofs given in [27] utilized Lin’s method, while we shall transform our system near the periodic orbit into Fenichel’s normal form and use reversible Shilnikov variables for the resulting system.

The rest of the paper is organised as follows. In §2, we state our hypotheses and the main results for systems in $\mathbb{R}^4$. We introduce reversible Shilnikov variables in §3, and the proofs of the main results for symmetric...
and asymmetric pulses are given in §4 and §5, respectively. Various extensions, for instance to isolas and to higher-dimensional systems including 2D patterns, are outlined and proved in §6. We end in §7 with a discussion of the Swift–Hohenberg equation and the stability of symmetric and asymmetric localised states with respect to the underlying partial differential equation (PDE).

2 Main results

We consider ordinary differential equations

\[ u_x = f(u, \mu), \quad u \in \mathbb{R}^4, \quad \mu \in \mathbb{R}, \]  

where \( f \) is assumed to be a smooth function. We are interested in reversible systems that conserve an additional real-valued quantity which we refer to as the energy. More precisely, we assume the following.

**Hypothesis 1** There exists a linear map \( R : \mathbb{R}^4 \to \mathbb{R}^4 \) with \( R^2 = 1 \) and \( \dim \text{Fix } R = 2 \) so that \( f(Ru, \mu) = -Rf(u, \mu) \) for all \((u, \mu)\).

From the reversibility hypothesis 1, we can infer that \( Ru(-x) \) satisfies (2.1) whenever \( u(x) \) does. We say that a solution \( u(x) \) is symmetric or reversible if \( u(0) \in \text{Fix } R \), which implies that \( u(x) = Ru(-x) \) for all \( x \). Next, we assume that (2.1) has a first integral \( \mathcal{H} \) that is invariant under the reverser \( R \).

**Hypothesis 2** There exists a smooth function \( \mathcal{H} : \mathbb{R}^4 \times \mathbb{R} \to \mathbb{R}^4 \) with \( \langle \nabla_u \mathcal{H}(u, \mu), f(u, \mu) \rangle = 0 \) and \( \mathcal{H}(Ru, \mu) = \mathcal{H}(u, \mu) \) for all \((u, \mu)\). We normalise \( \mathcal{H} \) so that \( \mathcal{H}(0, \mu) = 0 \) for all \( \mu \).

We can now state our assumptions on the specific solutions of (2.1) we are interested in.

**Hypothesis 3** The origin \( u = 0 \) is a hyperbolic equilibrium of (2.1). More precisely, we assume that \( f(0, \mu) = 0 \) for all \( \mu \) and that \( f_u(0, \mu) \) has two eigenvalues with strictly negative real part and another two eigenvalues with strictly positive real part.

Hypothesis 1 implies that the spectrum of \( f_u(0, \mu) \) is invariant under multiplication by \(-1\). Next, we require that (2.1) has, for each \( \mu \) in an appropriate interval, a symmetric periodic orbit in the zero energy level set. Throughout, we reserve the term “periodic orbit” for solutions with nonzero minimal period.

**Hypothesis 4** We assume that there is a closed interval \( J \subset \mathbb{R} \) with nonempty interior \( \dot{J} \) so that (2.1) has, for each \( \mu \in J \), a periodic orbit \( \gamma(x, \mu) \) with minimal period \( \ell(\mu) \) which satisfy:

(i) The family \( \gamma(x, \mu) \) depends smoothly on \( \mu \in J \).

(ii) \( \gamma(x, \mu) \) is symmetric with \( \gamma(0, \mu) \in \text{Fix } R \) for all \( \mu \in J \).

(iii) \( \gamma(x, \mu) \) has zero energy: for each \( \mu \in J \), we have \( \mathcal{H}(\gamma(x, \mu), \mu) = 0 \) and \( \mathcal{H}_u(\gamma(x, \mu), \mu) \neq 0 \) for one, and hence all, \( x \).

(iv) \( \gamma(x, \mu) \) has two positive nontrivial Floquet multipliers\(^2\) \( e^{\pm 2\pi \alpha(\mu)} \) with \( \alpha(\mu) > 0 \) for all \( \mu \in J \).

\(^2\)The particular form of the Floquet exponents facilitates a rescaling of the period of \( \gamma \) to \( 2\pi \).
Rescaling time if necessary, we can, without loss of generality, assume that the minimal periods \( \ell(\mu) \) are all equal to \( 2\pi \). In particular, \[38\] implies that \( \gamma(\pi, \mu) \in \text{Fix}(R) \). We will always identify \( S^1 = [0, 2\pi]/\sim \).

Some of our results change if \( (2.1) \) has additional symmetries. To account for such systems, we shall sometimes assume that \( (2.1) \) respects the reflection symmetry \( \kappa : u \mapsto -u \). To formulate this hypothesis, we need to introduce the variational equation

\[ v_x = f_u(\gamma(x, \mu), \mu) v \]

about \( \gamma(x, \mu) \) which, due to Hypothesis 4(iv), has two nontrivial solutions \( v^a \) and \( v^u \) of the form

\[ v^a(x) = e^{-\alpha(\mu)x} p^a(x, \mu), \quad v^u(x) = e^{\alpha(\mu)x} p^u(x, \mu), \quad (2.2) \]

where \( p^a(x, \mu) \) are \( 2\pi \)-periodic in \( x \). We shall normalize these functions by requiring that \( p^u(x, \mu) = R p^a(-x, \mu) \) for all \( x \) and \( \mu \in J \).

**Hypothesis 5** The nonlinearity \( f \) is odd in \( u \) so that \( f(-u, \mu) = -f(u, \mu) \) for all \( (u, \mu) \). Furthermore, we assume that the orbit \( \gamma(\cdot, \mu) \) is invariant under \( \kappa : u \mapsto -u \) and that \( p^a(\pi, \mu) = \kappa p^u(0, \mu) \) for all \( \mu \in J \).

Hypothesis 5 implies that \( \kappa R = -R \) is a second reverser of \( (2.1) \). Furthermore, the first part of Hypothesis 5 implies that the flow \( \Phi_\kappa(u) \) of \( (2.1) \) satisfies \( \kappa \Phi_\kappa(u) = \Phi_\kappa(\kappa u) \), and therefore \( D\Phi_\kappa(u) = D\Phi_\kappa(-u) \), for all \( x \) and \( u \). If the orbit \( \gamma(\cdot, \mu) \) is invariant under \( \kappa \), then necessarily \( \kappa \gamma(x, \mu) = \gamma(x + \pi, \mu) \) for all \( x \). In particular, we have \( D\Phi_\kappa(\gamma(0, \mu)) = D\Phi_\kappa(\gamma(\pi, \mu)) \), which implies that \( p^a(\pi, \mu) = \pm p^u(0, \mu) \). The second part of Hypothesis 5 assumes that the negative sign occurs in our system.

**Lemma 1** Assume that Hypotheses 1-4 are met, then there exist a \( \delta > 0 \), a smooth reversible change of coordinates near \( \gamma(\cdot, \mu) \), and smooth real-valued functions \( h^c \), \( h^+_j \) and \( h^-_j \) for \( j = 1, 2 \) so that \( (2.1) \) restricted to the zero energy level set is for all \( \mu \in J \) of the form

\[
\begin{align*}
v^c_x &= 1 + h^c(v, \mu) v^a + v^u \\
v^a_x &= -[\alpha(\mu) + h^+_1(v, \mu) v^a + h^+_2(v, \mu) v^u] v^a \\
v^u_x &= [\alpha(\mu) + h^-_1(v, \mu) v^a + h^-_2(v, \mu) v^u] v^u
\end{align*}
\]

where \( v = (v^c, v^a, v^u) \in \mathcal{V} := S^1 \times I \times I \) for \( I = [-\delta, \delta] \). The reverser \( R \) and the symmetry \( \kappa \) (if Hypothesis 5 is met) act on \( v \) via

\[ R(v^c, v^a, v^u) = (-v^c, v^a, v^u), \quad \kappa(v^c, v^a, v^u) = (v^c + \pi, v^a, v^u). \quad (2.4) \]

We shall prove Lemma 1 in §3. Note that \( v^u = 0 \) corresponds to the unstable manifold of the periodic orbit \( \gamma(\cdot, \mu) \) and its strong unstable fibers \( W^u(\gamma(x_0, \mu), \mu) \) are given by further fixing \( \varphi = x_0 \). We define the sections

\[ \Sigma_{in} = S^1 \times \{ v^u = \delta \} \times I, \quad \Sigma_{out} = S^1 \times I \times \{ v^a = \delta \}. \]

We are interested in connecting orbits between the periodic orbits \( \gamma(x, \mu) \) and the equilibrium \( u = 0 \), and therefore define

\[ \Gamma := \{ (\varphi, \mu) \in S^1 \times J; W^s(0, \mu) \cap \cap W^u(\gamma(\varphi, \mu), \mu) \cap \Sigma_{out} \neq \emptyset \}. \quad (2.5) \]

As in the introduction, we now assume that \( \Gamma \) is a graph and refer to §6.1 for extensions to isolas of fronts.

**Hypothesis 6** The set \( \Gamma \) is the graph of a smooth function \( z : S^1 \to J \). Furthermore, there exist an open neighbourhood \( U_\Gamma \) of \( \Gamma \) in \( S^1 \times J \), positive constants \( \epsilon, b > 0 \), and a smooth function \( g : U_\Gamma \to I \) so that

\[ \{ v \in W^s(0, \mu) \cap \Sigma_{out}; |v^a| < \epsilon, (\varphi, \mu) \in U_\Gamma \} = \{ (\varphi, g(\varphi, \mu), \delta) ; (\varphi, \mu) \in U_\Gamma \} \]

and \( |g(\varphi, \mu)| \geq b > 0 \) for all \( (\varphi, \mu) \in U_\Gamma \).
Hypothesis 6 encodes the bifurcation diagram of fronts or heteroclinic orbits that connect the periodic orbits \( \gamma \) to \( u = 0 \) through the intersections of the strong unstable fibers \( W^{uu}(\gamma(x,\mu),\mu) \) and the stable manifold \( W^s(0,\mu) \). Nondegenerate maxima and minima of \( z \) correspond to saddle-node bifurcations of fronts, that is, to quadratic tangencies of \( W^s(0,\mu) \) and \( W^{uu}(\gamma(x,\mu),\mu) \) at which heteroclinic connections disappear. For our results on asymmetric pulses, we need to assume that the only critical points of \( z \) are nondegenerate maxima or minima which occur at different values of \( z \).

**Hypothesis 7** We assume that \( z'(z) = 0 \) implies \( z''(z) \neq 0 \). Furthermore, we assume that \( z(z_1) = z(z_2) \) and \( z'(z_1) = z'(z_2) = 0 \) implies \( z_1 = z_2 \) modulo \( 2\pi \) (modulo \( \pi \) if Hypothesis 5 is met).

We define

\[
\Theta := \begin{cases} 
\{0, \pi/2\} & \text{if Hypothesis 5 is met} \\
\{0, \pi\} & \text{otherwise.}
\end{cases}
\]

We can now state our two main results.

**Theorem 1 (Symmetric 1-pulses)** Assume that Hypotheses 1-4 and 6-7 are met, then there are constants \( L_\ast \gg 1 \) and \( \eta > 0 \) so that the following is true: for each \( L > L_\ast \), (2.1) has a symmetric 1-homoclinic orbit \( u(x) \) for \( \mu \in J \) that spends time \( 2L \) in the neighbourhood \( V \) of \( \gamma(\cdot,\mu) \) if and only if

\[
\mu = \mu_\ast(L, \varphi_0) = z(\varphi_0 + L) + O(e^{-\eta L})
\]

for an appropriate \( \varphi_0 \in \Theta \), and \( u(0) \) lies near \( \gamma(\varphi,\mu) \) in \( \text{Fix}(R) \) for \( \varphi_0 \in \{0, \pi\} \).

If Hypothesis 5 is met, too, then \( z \) is \( \pi \)-periodic, the above statements are true for \( \varphi_0 \in \{0, \pi/2\} \), and \( u(0) \) lies near \( \gamma(\pi/2,\mu) \) in \( \text{Fix}(-R) \) when \( \varphi_0 = \pi/2 \). Furthermore, symmetric 1-pulses come in pairs: if \( u \) is a 1-pulse corresponding to \( \varphi_0 \), then \( -u \) is a 1-pulse corresponding to \( \varphi_0 + \pi \).

Lemma 4 contains precise expansions of \( \mu_\ast \) and the associated fold bifurcations that sharpen the estimate given in the preceding theorem.

**Theorem 2 (Asymmetric 1-pulses)** Assume that Hypotheses 1-4 and 6-7 are met. We denote by \( k \) the number of maxima of \( z \) and by \( \mathcal{S}_{\varphi_0} \) with \( \varphi_0 \in \Theta \) the two branches of symmetric 1-pulses described in Theorem 1. There are then an \( \eta > 0 \), natural numbers \( n_\ast \gg 1 \) and \( k \geq 0 \), and smooth curves

\[
a_i, a_i(n, \varphi_0), \tilde{a}_j(n, \varphi_0) : [0, 1] \rightarrow \mathbb{R}^2, \quad 1 \leq i \leq k, \quad 1 \leq j \leq \tilde{k}
\]

defined for \( n > n_\ast \) and \( \varphi_0 \in \Theta \) so that the following is true for \( n > n_\ast \) and \( \varphi_0 \in \Theta \):

(i) We have

\[
a_i(s, n, \varphi_0) = a_i(s) + (0, \varphi_0 + 2\pi n) + O(e^{-\eta n}), \quad s \in [0, 1], \quad 1 \leq i \leq k.
\]

For \( 1 \leq i < k \), the curves \( a_i(s, n, \varphi_0) \) begin for \( s = 0 \) in \( \mathcal{S}_{\varphi_0} \) near a non-global maximum\(^3 \) of \( z \) and terminate for \( s = 1 \) in \( \mathcal{S}_{\varphi_0} \) near a non-global minimum of \( z \). The curves \( a_k(s, n, \varphi_0) \) begin for \( s = 0 \) in \( \mathcal{S}_{\varphi_0} \) near the global maximum of \( z \) and terminate for \( s = 1 \) in \( \mathcal{S}_{\varphi_0 + \pi} \) near the global minimum of \( z \).

(ii) We have

\[
\tilde{a}_j(s, n, \varphi_0) = \tilde{a}_j(s) + (0, \varphi_0 + 2\pi n) + O(e^{-\eta n}), \quad \tilde{a}_j(0, n, \varphi_0) = \tilde{a}_j(1, n, \varphi_0), \quad s \in [0, 1], \quad 1 \leq j \leq \tilde{k}.
\]

\(^3\)That is, near a point \((\mu_\ast, L_\ast)\) on \( \mathcal{S}_{\varphi_0} \) given by (2.6) where \( z(\varphi_0 + L) \) has a non-global maximum at \( L = L_\ast \).
(iii) Equation (2.1) has an asymmetric 1-homoclinic orbit $u(x)$ for $\mu \in J$ that spends time $2L$ in the neighborhood $V$ of $\gamma(\cdot, \mu)$ if and only if $(\mu, L)$ lies on one of the curves $a_i(s, n, \varphi_0)$ with $s \neq 0, 1$ or on $\tilde{a}_j(s, n, \varphi_0)$ for $s \in [0, 1]$.

(iv) The start and end points of the curves $a_i(s, n, \varphi_0)$ for $s = 0, 1$ correspond to pitchfork bifurcations of the symmetric 1-pulses described in Theorem 1. No other pitchfork bifurcations occur on $S_{\varphi_0}$.

(v) The curves $\tilde{a}_j(\cdot, n, \varphi_0)$ correspond to isolas of asymmetric solutions. Furthermore, the curves $a_j(s)$ and $\tilde{a}_j(s)$ are obtained from solutions $(L, \varphi)$ of (1.5) upon setting $\mu = z(L + \varphi)$.

Finally, if Hypothesis 5 is met, then the above statements remain true provided we replace $S_{\varphi_0 + \pi}$ in (i) by $S_{\varphi_0 + \pi/2}$. Furthermore, if $u$ is an asymmetric 1-pulse, so are $Ru$, $-u$, and $-R\mu$.

We refer to Lemma 5 for expansions of the location of the pitchfork bifurcations that corroborate the summary given in Figure 9. Lemma 8 shows that isolas are only possible when $k > 2$, and an example of a function $z$ that leads to isolas is given in Figure 8.

3 Reversible Shilnikov variables near periodic orbits

We begin with the proof of Lemma 1.

Proof of Lemma 1 Hypotheses 4(iii)-(iv) imply that the zero energy level set near $\gamma(\cdot, \mu)$ is diffeomorphic to $S^1 \times I \times I$. For instance, we may use the coordinates

$$v = (v^r, v^s, v^a) \in S^1 \times I \times I \mapsto \gamma(v^r, \mu) + v^r \hat{p}^r(v^r, \mu) + v^a \hat{p}^a(v^r, \mu) + h.o.t., \quad (3.1)$$

which utilize the Floquet solutions from (2.2), to parametrize a neighbourhood of the periodic orbits in the zero energy level set. Straightening out the invariant stable and unstable manifolds of the periodic orbits and their stable and unstable fibers transforms (2.1) into (2.3) in this neighbourhood [9]. It is not hard to see that this coordinate transformation, if done correctly, preserves the reverser $R$ and the symmetry $\kappa$ (if Hypothesis 5 is met). Furthermore, using (3.1) together with $\gamma(v^r, \mu) \in \text{Fix}(R)$ for $v^c = 0, \pi$, the normalization of $p^r(x, \mu)$ after (2.2), and Hypothesis 5, we find that the action of $R$ and $\kappa$ on the $v$ are given by (2.4) as claimed.

We now solve (2.3) near the periodic orbits.
Lemma 2 There exist positive constants $L_0$ and $\eta$ so that the following is true for all $L > L_0$ and $\varphi \in S^1$: there is a unique solution $v(x)$, also referred to as $v(x, \varphi)$, of (2.3), defined for $x \in [-L, L]$, so that

$$v(-L) \in \Sigma_{in}, \quad v(L) \in \Sigma_{out}, \quad v'(0) = \varphi, \quad v(x) \in V \quad \forall x \in [-L, L].$$

Furthermore, we have

$$v(-L) = \left(\varphi - L + O(e^{-\eta L}), \delta, \delta e^{-\alpha(\mu) L}(1 + O(e^{-\eta L}))\right),$$

$$v(L) = \left(\varphi + L + O(e^{-\eta L}), \delta e^{-\alpha(\mu) L}(1 + O(e^{-\eta L})), \delta\right)$$

and

$$v(0) = \left(\varphi, \delta e^{-\alpha(\mu) L}(1 + O(e^{-\eta L})), \delta e^{-\alpha(\mu) L}(1 + O(e^{-\eta L}))\right).$$

(3.2)

The solution $v(x)$ is smooth in $(\varphi, \mu, L)$, and the error estimates in (3.2) can be differentiated. Furthermore, we have

$$v(x, -\varphi) = Rv(-x, \varphi), \quad \varphi \in S^1, \ |x| \leq L$$

(3.3)

and, if Hypothesis 5 holds,

$$v(x, \pi + \varphi) = \kappa v(x, \varphi), \quad \varphi \in S^1, \ |x| \leq L$$

(3.4)

where $R$ and $\kappa$ act according to (2.4). In particular, the solution $v(x, \varphi)$ is $R$-reversible, with $v(0) \in \text{Fix}(R)$, if and only if $\varphi = 0, \pi$, and $\kappa R$-reversible, with $v(0) \in \text{Fix}(\kappa R)$, if and only if $\varphi = \pi/2, 3\pi/2$.

Proof. Existence and uniqueness follow, for instance, from [19, Theorem 4 and §5] or [37, Theorem 2.1]. The expansion for $\varphi(x)$ follows from [37, Theorem 2.1], while the expansions for $v'(x)$ and $v''(x)$ follow then as in [19, §5] or [32]. The claims about reversibility follow from the expression for $v'(0)$ together with the uniqueness of the solution considered in the lemma. Note that (2.4) implies that $R \Sigma_{in} = \Sigma_{out}$ and that the elements on the periodic orbits fixed by $R$ and $\kappa R$ are $\varphi = 0, \pi$ and $\varphi = \pi/2, 3\pi/2$, respectively.

In the proofs given in §4 and §5, we shall denote different positive exponential rates $\eta$ that appear in estimates of remainder terms by the same constant, as they are bounded from below by the same positive constant.

Next, we consider the intersection of $W^u(0, \mu)$ with $\Sigma_{in}$.

Lemma 3 We have $v = (\varphi, \delta, v') \in W^u(0, \mu)$ if and only if $v = (-\varphi, v^u, \delta) \in W^s(0, \mu)$. If Hypothesis 5 holds, then $v \in W^s(0, \mu) \cap \Sigma_{out}$ if and only if $\kappa v \in W^s(0, \mu) \cap \Sigma_{out}$. In particular, $g(\varphi, \mu)$ and $z(\varphi)$ are then $\pi$-periodic in $\varphi$ for all $(\varphi, \mu) \in U_\Gamma$.

Proof. The statements follow from reversibility and equivariance of (2.1) and the coordinate transformation in Lemma 1 together with the representation of $R$ in (2.4).

In particular, both $g$ and $z$ are $\pi$-periodic in $\varphi$ if Hypothesis 5 is met, which explains why we imposed the conditions in Hypothesis 7 only modulo $\pi$ in this case. We record also that, on account of Hypothesis 6, we have

$$g(\varphi, z(\varphi)) \equiv 0 \quad \forall \varphi, \quad g(\varphi, \mu) = 0 \quad \text{if and only if} \quad \mu = z(\varphi).$$

(3.5)

Using Lemma 2 and 3, we can now construct symmetric and asymmetric homoclinic orbits to $u = 0$ that pass near the periodic orbits $\gamma(\cdot, \mu)$. 

11
4 Symmetric 1-pulses

In this section, we are interested in finding reversible homoclinic orbits $v(x)$ to $u = 0$ that spend a long time near the periodic orbit $\gamma(\cdot, \mu)$. We focus on reversible 1-pulses which satisfy by definition

\[
v(x) \in V \iff x \in [-L, L]
\]
\[
v(L) \in W^s(0, \mu) \cap \Sigma_{\text{out}}
\]
\[
v(0) \in \text{Fix}(R) \text{ or } v(0) \in \text{Fix}(\kappa R)
\]

for sufficiently large $L \gg 1$. Lemma 2 implies that (4.1) is met if and only if

\[
v(L) = (v^c(L, \varphi_0, \mu), v^s(L, \varphi_0, \mu), \delta) \in W^s(0, \mu),
\]

where $\varphi_0 \in \Theta$ is fixed. For $L \gg 1$, Hypothesis 6 shows that this equation is met if and only if

\[
(v^c(L, \varphi_0, \mu), v^s(L, \varphi_0, \mu)) = (\varphi, g(\varphi, \mu))
\]

for some $(\varphi, \mu) \in U_\Gamma$ or, equivalently,

\[
g(v^c(L, \varphi_0, \mu), \mu) = v^s(L, \varphi_0, \mu).
\]

We have

\[
v^c(L, \varphi_0, \mu) \overset{(3.2)}{=} \varphi_0 + L + O(e^{-\eta L}), \quad \varphi_0 \in \Theta \text{ fixed}
\]

and define

\[
\ell := v^c(L, \varphi_0, \mu) - \varphi_0 = L + O(e^{-\eta L}),
\]

where we treat $\ell$ as an element of the universal cover $\mathbb{R}$ of $S^1$. Thus, $\ell$ measures the phase in $\Sigma_{\text{out}}$, while $2L$ is the time spent by the solution $v(x)$ near the periodic orbit $\gamma(\cdot, \mu)$. We can solve (4.3) uniquely for $L \gg 1$ to get

\[
L = L_*(\ell, \mu) = \ell + O(e^{-\eta \ell}), \quad \ell \gg 1, \quad \mu \in J.
\]

Substituting (4.4) into (4.2), we obtain the equation

\[
g(\varphi_0 + \ell, \mu) = v^s(L_*(\ell, \mu), \varphi_0, \mu) \overset{(3.2)}{=} \delta e^{-2\alpha(\mu)\ell}(1 + O(e^{-\eta \ell})).
\]

It suffices now to solve (4.5) for $\ell \gg 1$ and $\mu \in J$, where $\varphi_0 \in \Theta$ is fixed.

**Lemma 4** Fix $\varphi_0 \in \Theta$, then there exists an $\ell_0 \gg 1$ so that (4.5) has a unique solution $\mu = \mu_*(\ell, \varphi_0) \in J$ for each $\ell > \ell_0$. Furthermore,

\[
\mu_*(\ell, \varphi_0) = z(\varphi_0 + \ell) + \frac{\delta e^{-2\alpha(z(\varphi_0 + \ell))}}{g\mu(\varphi_0 + \ell, z(\varphi_0 + \ell))}(1 + O(e^{-\eta \ell}))
\]

for all $\ell > \ell_0$. Near each $\ell_*$ with $z'(\varphi_0 + \ell_*) = 0$, there is a unique fold bifurcation of the symmetric 1-pulses: the fold is nondegenerate and occurs at

\[
\ell = \ell_* + \left(\frac{2\alpha}{g_0^{2\eta}} + \frac{g_\varphi \mu}{g_0^{3\eta}}\right) \delta e^{-2\alpha \ell_*}(1 + O(e^{-\eta \ell_*}))
\]

where all terms are evaluated at $(\varphi_0 + \ell_*, z(\varphi_0 + \ell_*))$. No other fold bifurcations occur.
Proof. We set \( u := z(\varphi_0 + \ell) + \tilde{\mu} \) so that (4.5) becomes

\[
g(\varphi_0 + \ell, z(\varphi_0 + \ell) + \tilde{\mu}) = \delta e^{-2\alpha(z(\varphi_0 + \ell) + \tilde{\mu})f(1 + O(e^{-\eta f}))}.
\]

On account of (3.5), \( g(\varphi, \mu) = 0 \) if and only if \( \mu = z(\varphi) \). Thus, the above equation can be written as

\[
\tilde{\mu} + O(\mu^2) = \frac{\delta}{g_{\mu}(\varphi_0 + \ell) + z(\varphi_0 + \ell)} e^{-2\alpha(z(\varphi_0 + \ell) + \tilde{\mu})f(1 + O(e^{-\eta f}))}
\]

where we used that \( |g_{\mu}| \geq b > 0 \) in \( \Gamma \) by Hypothesis 6. Using the implicit function theorem, we readily obtain (4.6). Fold bifurcations occur precisely when \( H \) of Hypothesis 7.

Theorem 1 is a direct consequence of Lemma 4.

5 Asymmetric 1-pulses

We now consider 1-pulses that are not necessarily symmetric. Throughout this section, we assume that Hypotheses 6-7 are met. We first consider the case where we do not have an additional \( \mathbb{Z}_2 \)-symmetry, that is, we do not assume Hypothesis 5, and shall comment in \$5.4$ below on the case where Hypothesis 5 is met. Arbitrary 1-pulses correspond to solutions \( v(x) \) with

\[
\begin{align*}
v(x) & \in V \iff x \in [-L, L] \\
v(-L) & \in W^u(0, \mu) \cap \Sigma_{in} \\
v(L) & \in W^s(0, \mu) \cap \Sigma_{out}
\end{align*}
\]

for sufficiently large \( L \gg 1 \). Using Lemmas 2 and 3, we see that (5.1) is equivalent to \( G(L, \varphi, \mu) = 0 \) where

\[
G(L, \varphi, \mu) = \begin{pmatrix} g(v^c(L, \varphi, \mu), \mu) - v^s(L, \varphi, \mu) \\ g(-v^c(-L, \varphi, \mu), \mu) - v^s(-L, \varphi, \mu) \end{pmatrix} = \begin{pmatrix} g(v^c(L, \varphi, \mu), \mu) - v^s(L, \varphi, \mu) \\ g(v^c(L, -\varphi, \mu), \mu) - v^s(L, -\varphi, \mu) \end{pmatrix}
\]

and \( \varphi \in S^1 \) is arbitrary.

5.1 Pitchfork bifurcations from symmetric 1-pulses

We begin by investigating bifurcations from the symmetric 1-pulses. We set

\[
\tilde{G}(L, \varphi, \mu) := \begin{pmatrix} \tilde{G}_1(L, \varphi, \mu) + \tilde{G}_2(L, \varphi, \mu) \\ \tilde{G}_1(L, \varphi, \mu) - \tilde{G}_2(L, \varphi, \mu) \end{pmatrix} \equiv \begin{pmatrix} \tilde{G}_1(L, \varphi, \mu) + \tilde{G}_1(L, -\varphi, \mu) \\ \tilde{G}_1(L, \varphi, \mu) - \tilde{G}_1(L, -\varphi, \mu) \end{pmatrix}
\]

where \( (\tilde{G}_1, \tilde{G}_2) \) are the components of \( \tilde{G} \) defined in (5.2). Note that \( \tilde{G}_2(L, \varphi_0, \mu) \equiv 0 \) for \( \varphi_0 \in \Theta \) and that the symmetric 1-pulses found in \$4$ correspond to solutions of \( \tilde{G}_1(L, \varphi_0, \mu) = 0 \) with \( \varphi_0 \in \Theta \). Pitchfork bifurcations of symmetric 1-pulses can occur only when

\[
\tilde{G}_1(L, \varphi_0, \mu) = 0, \quad \partial_2 \tilde{G}_2(L, \varphi_0, \mu) = 0
\]

for some \( \varphi_0 \in \Theta \). Inspecting (5.3), we see that these equations are equivalent to

\[
G_1(L, \varphi_0, \mu) = 0, \quad \partial_2 G_1(L, \varphi_0, \mu) = 0,
\]
that is, to
\[ g(v^c(L, \varphi_0, \mu), \mu) - v^s(L, \varphi_0, \mu) = 0, \quad \partial_\varphi [g(v^c(L, \varphi, \mu), \mu) - v^s(L, \varphi, \mu)]|_{\varphi=\varphi_0} = 0 \] (5.4)
for some \( \varphi_0 \in \Theta \). The following lemma gives the location of pitchfork bifurcations and asserts that these lead to locally unique branches of asymmetric 1-pulses that are parametrized by \( \varphi \) near \( \varphi_0 \).

**Lemma 5** Fix \( \varphi_0 \in \Theta \), then there are numbers \( \rho > 0 \) and \( \ell_0 \gg 1 \) with the following property. Near each \( \ell_\ast > \ell_0 \) with \( z'(\varphi_0 + \ell_\ast) = 0 \), there is a unique pitchfork bifurcation of the symmetric 1-pulses, occurring for \( \mu = \mu_\ast(\ell, \varphi_0) \) as in (4.6) and
\[ \ell = \ell_\ast + \frac{g_{\varphi_0}}{g_\mu^2} \delta e^{-2\alpha \ell_\ast}(1 + O(e^{-\eta \ell_\ast})) \] (5.5)
where all terms are evaluated at \( (\varphi_0 + \ell_\ast, z(\varphi_0 + \ell_\ast)) \) and \( \ell \) is related to \( L \) via (4.3). At these pitchfork bifurcations, a locally unique branch of asymmetric 1-pulses bifurcates, and this branch is parametrized smoothly by \( \varphi \) for \( \varphi \in U_\rho(\varphi_0) \). No other pitchfork bifurcations occur.

**Proof.** Pitchfork bifurcations are captured by (5.4). We have already solved the first equation of (5.4) in Lemma 4, obtaining
\[ \mu = \mu_\ast(\ell, \varphi_0) = z(\varphi_0 + \ell) + \frac{\delta e^{-2\alpha (z(\varphi_0 + \ell)) \ell}}{g_\mu(\varphi_0 + \ell, z(\varphi_0 + \ell))}(1 + O(e^{-\eta \ell})) =: z(\varphi_0 + \ell) + \tilde{\mu}(\ell) \] (5.6)
with \( \ell \) given in (4.3). The remaining second equation of (5.4),
\[ \partial_\varphi G_1(L, \varphi_0, \mu) = g_\varphi(v^c(L, \varphi_0, \mu), \mu)v^c_\varphi(L, \varphi_0, \mu) - v^s_\varphi(L, \varphi_0, \mu) = 0, \]
is evaluated at \( (L, \mu) = (L_\ast(\ell, \mu_\ast(\ell, \varphi_0)), \mu_\ast(\ell, \varphi_0)) \), and equations (3.2), (4.3), and (5.6) show that it is given by
\[ g_\varphi(\varphi_0 + \ell, z(\varphi_0 + \ell) + \tilde{\mu}(\ell))(1 + O(e^{-\eta \ell})) - O(e^{-2\alpha (z(\varphi_0 + \ell)) + \eta \ell}) = 0. \] (5.7)
We expand as follows:
\[ g_\varphi(\varphi_0 + \ell, z(\varphi_0 + \ell) + \tilde{\mu}(\ell)) = g_\varphi(\varphi_0 + \ell, z(\varphi_0 + \ell)) + g_{\varphi \mu}(\varphi_0 + \ell, z(\varphi_0 + \ell))\tilde{\mu}(\ell) + O(\tilde{\mu}(\ell)^2) \]
(\( \equiv \)
\[ -g_\mu(\varphi_0 + \ell, z(\varphi_0 + \ell))z'(\varphi_0 + \ell) + g_{\varphi \mu}(\varphi_0 + \ell, z(\varphi_0 + \ell))\tilde{\mu}(\ell) + O(\tilde{\mu}(\ell)^2) \]
so that (5.7) becomes
\[ -g_\mu(\varphi_0 + \ell, z(\varphi_0 + \ell))z'(\varphi_0 + \ell) + g_{\varphi \mu}(\varphi_0 + \ell, z(\varphi_0 + \ell))\tilde{\mu}(\ell) + O(e^{-2\alpha (z(\varphi_0 + \ell)) + \eta \ell}) = 0. \] (5.8)
Since \( g_\mu \) is bounded away from zero by Hypothesis 6 and \( \tilde{\mu}(\ell) \to 0 \) as \( \ell \to \infty \), the above equation can be satisfied only near points \( \ell_\ast \) with \( z'(\varphi_0 + \ell_\ast) = 0 \). Setting \( \ell = \ell_\ast + \hat{\ell} \), expanding (5.8) around \( \ell_\ast \), and using that \( z''(\varphi_0 + \ell_\ast) \neq 0 \) by Hypothesis 7 gives (5.5).

To prove that a locally unique branch bifurcates at each pitchfork bifurcation and that this branch can be parametrized by \( \varphi \) for \( \varphi \) near \( \varphi_0 \), it suffices to prove that
\[ \frac{d}{dL}[\partial_\varphi G_1(L, \varphi_0, \mu_\ast(L, \varphi_0))] \neq 0. \]
Using (4.4) and (5.8), we find
\[ \frac{d}{dL}[\partial_\varphi G_1(L, \varphi_0, \mu_\ast(L, \varphi_0))] = -g_\mu(\varphi_0 + \ell_\ast, z(\varphi_0 + \ell_\ast))z''(\varphi_0 + \ell_\ast) + O(e^{-\eta \ell_\ast}) \] (5.9)
at the solution \( \ell \) given in (5.5), and this expression is indeed bounded away from zero uniformly in \( \ell_\ast \) by Hypotheses 6 and 7. Finally, the neighbourhood \( U_\rho(\varphi_0) \) over which we can parametrize the branches in \( \varphi \) is uniform in \( \ell_\ast \), since the estimates in (5.8) and (5.9) and the dependence of (5.2) on \( \varphi \) are uniform in \( \ell_\ast \).
Remark 1 We briefly compare the locations of saddle-node and pitchfork bifurcations given in Lemmas 4 and 5, respectively. Near each \( \ell_\ast \) with \( z'(\varphi_0 + \ell_\ast) = 0 \), there is a unique saddle-node at

\[
\ell_{\text{sn}} = \ell_\ast + \left( \frac{2\alpha}{g_\mu z''} + \frac{g_\varphi \mu}{g_\mu z''} \right) \delta e^{-2\alpha \ell_\ast} \left( 1 + O(e^{-\eta \ell_\ast}) \right)
\]

and a unique pitchfork at

\[
\ell_{\text{pf}} = \ell_\ast + \frac{g_\varphi \mu}{g_\mu z''} \delta e^{-2\alpha \ell_\ast} \left( 1 + O(e^{-\eta \ell_\ast}) \right)
\]

where all terms are evaluated at \((\varphi_0 + \ell_\ast, z(\varphi_0 + \ell_\ast))\), see (4.7) and (5.5), respectively. Thus, we have

\[
\ell_{\text{sn}} = \ell_{\text{pf}} + \frac{2\alpha}{g_\mu z''} \delta e^{-2\alpha \ell_\ast} \left( 1 + O(e^{-\eta \ell_\ast}) \right),
\]

and the sign of the offset between these bifurcations is determined by the sign of the product \( g_\mu z'' \). The resulting scenarios are summarised in Figure 9.

5.2 Continuation of asymmetric 1-pulses

We showed in §5.1 that a locally unique branch of asymmetric 1-pulses bifurcates at each of the infinitely many pitchfork bifurcations identified in Lemma 5. In this section, we prove that each of these branches terminates at another pitchfork bifurcation, and we determine at which one.

To find asymmetric 1-pulses, we need to solve \( G(L, \varphi, \mu) = 0 \), where \( G \) is defined in (5.2). Lemma 2 implies that this equation is given by

\[

g(L + \varphi + O(e^{-\eta L}), \mu) = O(e^{-\eta L})
\]

\[

g(L - \varphi + O(e^{-\eta L}), \mu) = O(e^{-\eta L}),
\]

where \( \varphi \in S^1 \) is arbitrary, possibly after adjusting \( \eta \) to a smaller positive value.

For \( \varphi = \varphi_0 \in \Theta \), we recover the symmetric 1-pulses. For each \( \varphi_0 \in \Theta \) and near each \( L_\ast \) with \( z'(L_\ast + \varphi_0) = 0 \), Lemma 5 guarantees the existence of a locally unique branch of asymmetric solutions of (5.10) that are parametrized by \( \varphi \in U_\mu(\varphi_0) \). Furthermore, no other bifurcations to asymmetric 1-pulses occur. We are now interested in continuing the branches of asymmetric 1-pulses.

The first equation in (5.10) can be solved for \( \mu \) as in §4, and we find

\[
\mu = z(L + \varphi) + O(e^{-\eta L}).
\]

Using this expression, the second equation in (5.10) becomes

\[
g(L - \varphi + O(e^{-\eta L}), z(L + \varphi) + O(e^{-\eta L})) = O(e^{-\eta L}).
\]

We expect that the \( O(e^{-\eta L}) \) terms in (5.11) will be of higher order. Omitting these terms for a moment, we arrive at the equation

\[
g(L - \varphi, z(L + \varphi)) = 0
\]

which, on account of (3.5), has a solution \((L, \varphi)\) if and only if

\[
z(L - \varphi) = z(L + \varphi).
\]

Recall that \( z \) is \( 2\pi \)-periodic in its argument. Thus, for \( \varphi \in \Theta \), (5.13) is met for all \( L \), and these solutions correspond to symmetric 1-pulses. In Lemma 5, we found all solutions \((L, \varphi)\) of (5.10) near each \((L_\ast, \varphi_0)\) with \( \varphi_0 \in \Theta \) and \( z'(L_\ast + \varphi_0) = 0 \). Our goal is now to show that all other solutions of (5.10) can be found by solving (5.13).
Lemma 6 The function $G(L, \varphi) := g(L - \varphi, z(L + \varphi))$ is $2\pi$-periodic in both $\varphi$ and $L$, and we have $\nabla G(L, \varphi) \neq 0$ at each point $(L, \varphi)$ for which $z(L - \varphi) = z(L + \varphi)$ except when both $\varphi \in \Theta$ and $z'(L + \varphi) = 0$.

Proof. We have
\[
\nabla G(L, \varphi) = \begin{pmatrix} g_\varphi(L - \varphi, z(L + \varphi)) \\ g_\mu(L - \varphi, z(L + \varphi))z'(L + \varphi) \end{pmatrix}.
\]
Equation (5.13) means that the argument $(L - \varphi, z(L + \varphi)) = (L - \varphi, z(L - \varphi))$ lies in $\Gamma = \text{graph}(z)$. Hypothesis 6 then implies that $|g_\mu(L - \varphi, z(L + \varphi))| \geq b > 0$, and we can conclude that $\nabla G(\varphi, L) = 0$ if and only if both $z'(L + \varphi) = 0$ and $g_\varphi(L - \varphi, z(L + \varphi)) = g_\varphi(L - \varphi, z(L - \varphi)) = 0$. Using Hypothesis 6 and the definition of $g$, it is easy to see, via a Taylor expansion, that the second equation implies $z'(L - \varphi) = 0$. Thus, we have $z'(L \pm \varphi) = 0$ which, together with (5.13) and Hypothesis 7, implies $\varphi \in \Theta$ as claimed.

Lemma 6 will allow us below to solve (5.11) by the implicit function theorem near any solution $(L, \varphi)$ of (5.13) uniformly in $L \gg 1$, unless we have both $\varphi \in \Theta$ and $z'(L + \varphi) = 0$ in which case we already know all solutions by Lemma 5.

Thus, we initially focus on solving (5.13)
\[
Z(L, \varphi) := z(L + \varphi) - z(L - \varphi) = 0
\]
for $(L, \varphi)$. We shall see that the shape of the graph of $z$ determines the entire solution structure of (5.14) and refer to Figure 7 for an illustration. We set
\[
Q := S^1 \times [0, \pi], \quad \Lambda := \{(L, \varphi) \in Q; Z(L, \varphi) = 0\}, \quad \Lambda_{\text{bif}} := \{(L, \varphi) \in \partial Q; z'(L + \varphi) = 0\}.
\]
Equation (5.14) readily shows that $\Lambda_{\text{bif}} \subset \partial Q \subset \Lambda$. Furthermore, let
\[
M := \max_{L \in S^1} z(L) = z(L_M), \quad m := \min_{L \in S^1} z(L) = z(L_m)
\]
and note that $L_M$ and $L_m$ are unique in $S^1$ by Hypothesis 7. The following lemma contains the main result of this section, see also Figure 10.

Lemma 7 Assume that Hypotheses 6 and 7 are met, and let $k$ be the number of maxima of the function $z$.

(i) $\Lambda$ is the disjoint union of the one-dimensional manifolds $\partial Q$, $\Lambda_{\text{pt}}$ and $\Lambda_{\text{iso}}$: $\Lambda_{\text{pt}}$ is the disjoint union of the $2k$ distinct global solution branches described in (ii) and (iii) below, and $\Lambda_{\text{iso}}$ consists of finitely many isolas, each diffeomorphic to a circle, whose projection onto the $L$-axis is not onto.

(ii) Each element $(L, 0)$ and $(L, \pi)$ of $\Lambda_{\text{bif}}$ is a generic pitchfork bifurcation point, giving rise to a unique global branch of solutions of (5.14) in $Q$. These branches do not cross, and they begin and end in $\Lambda_{\text{bif}}$.

(iii) The branches starting at $(L_m, 0)$ and $(L_M, 0)$ terminate at $(L_M - \pi, \pi)$ and $(L_m - \pi, \pi)$, respectively. Each branch starting at $(L_1, 0)$, where $L_1 \neq L_M$ is a non-global maximum of $z$, terminates at a point $(L_2, 0)$ where $L_2 \neq L_m$ is a non-global minimum of $z$. Similarly, branches starting at $(L_1, \pi)$, where $L_1 \neq L_M$ is a non-global maximum of $z$, terminate at points $(L_2, \pi)$, where $L_2 \neq L_m$ is a non-global minimum of $z$.

Proof. We record that
\[
\nabla Z(L, \varphi) = \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} z'(L + \varphi) \\ z'(L - \varphi) \end{pmatrix}
\]
Figure 10: The set $\Lambda = Z^{-1}(0)$ corresponding to the function $z$ plotted in the right panel is shown. The set $\Lambda_{\text{bif}}$ is the union of the bullets located at $(L,0)$ and $(L,\pi)$ which correspond to maxima and minima of $z(L + \varphi)$. Also shown is the sign of $Z$ in the components of $Q \setminus \Lambda$ which is used to orient branches of $\Lambda$. In general, $\Lambda$ may also contain isolas.

Figure 11: The left panel shows representative solutions $L \pm \varphi$ of $z(L + \varphi) = z(L - \varphi)$ near the pitchfork bifurcation from $L = L_M$. The right panel illustrates that the solutions $L - \varphi$ and $L + \varphi$ are trapped in the intervals $(L_m, L_M)$ and $(L_M, L_m + 2\pi)$, respectively, with $0 < \varphi < \pi$ until either $L - \varphi = L_m$ and $L + \varphi = L_m + 2\pi$, or else $L = L_M$ and $\varphi = 0$.

and consider the planar Hamiltonian vector field

$$\begin{pmatrix} L_s \\ \varphi_s \end{pmatrix} = F(L, \varphi) := \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \nabla Z(L, \varphi)$$

(5.15)

associated with the energy $Z(L, \varphi)$, whose zero level set is $\Lambda$. Note that $\nabla Z$ vanishes at $(L, \varphi) \in \Lambda$ if and only if both $z'(L + \varphi) = 0$ and $z'(L - \varphi) = 0$. Hypothesis 7 asserts that this can only happen when $\varphi = 0$ or $\varphi = \pi$, that is, when $(L, \varphi) \in \Lambda_{\text{bif}}$. Thus, all equilibria of (5.15) in $\Lambda$ lie in $\Lambda_{\text{bif}}$, and they are hyperbolic saddles since

$$DF(L, \varphi_0) = 2 \begin{pmatrix} z''(L - \varphi_0) & 0 \\ 0 & -z''(L - \varphi_0) \end{pmatrix}, \quad \varphi_0 \in \{0, \pi\}$$

at each such equilibrium. Since $\nabla Z$ is non-zero everywhere else on $\Lambda$, it follows from Poincaré–Bendixson that $\Lambda \setminus \Lambda_{\text{bif}}$ is a one-dimensional manifold, namely the union of the heteroclinic orbits that emanate from and terminate at $\Lambda_{\text{bif}}$ and finitely many periodic orbits. This proves (ii). We have also shown all claims made in (i) except that we have not yet proved that isolas cannot extend from $L = 0$ to $L = 2\pi$. This will follow at once from (iii) since the branches that connect global extrema serve as a barrier; see Figure 10.

If $(L, \varphi_0) \in \Lambda_{\text{bif}}$ corresponds to a maximum of $z$, then the eigenvalue belonging to the eigenvector $(0,1)$ is positive, and it is the unstable manifold of the equilibrium that evolves into the interior $\hat{Q}$ of $Q$. Similarly, at minima of $z$, it is the stable manifold that points into $\hat{Q}$. Thus, orbits emanating from maxima of $z$ must terminate at minima of $z$, and vice versa; see also Figure 10. This proves part of (iii).
Consider now the pitchfork branch emerging from \((L, \varphi) = (L_M, 0)\) where \(z(L_M)\) is the uniquely assumed maximal value of \(z\). We wish to prove that this branch terminates at \((L, \varphi) = (L_m, \pm \pi, \pi)\). To prove this assertion, note that \(L - \varphi < L_M < L + \varphi\) near the bifurcation point, see Figure 11. Thus, inspecting Figure 11, we see that there are precisely two possibilities: first, \(L - \varphi\) or \(L + \varphi\) reach \(L_m\) or \(L_m + 2\pi\) at some point, in which case both must assume these values as \(L_m\) is the only value modulo \(2\pi\) for which \(z(L) = z(L_m)\); in this case, \(\varphi = \pi\), and the branch terminates at \(L = L_m \pm \pi\). The second option is that \(L - \varphi\) or \(L + \varphi\) reach \(L_M\), in which case they both must assume this value together; this is ruled out though by the vector-field interpretation discussed above. Thus, only the first possibility can occur which proves the assertion.

Next, we consider a branch emerging from \((L, \varphi) = (L_s, 0)\) where \(L_s\) corresponds to a non-global maximum or minimum of \(z\). Without loss of generality, we may assume that \(L_s \in (L_m, L_M)\). We conclude that \(L - \varphi\) and \(L + \varphi\) can never leave the interval \((L_m, L_M)\) along the branch for \(\varphi > 0\) since \(m < z(L - \varphi) = z(L + \varphi) < M\) for all \(L \pm \varphi \in (L_m, L_M)\), see also Figure 11. Thus, the branch must terminate at \(\varphi = 0\) as claimed. This finishes the proof of (iii).

The results established in Lemma 7 are valid for (5.13) and therefore also for (5.12). We need to lift these results to the full problem (5.11) given by

\[
g(L - \varphi + O(e^{-nL}), z(L + \varphi) + O(e^{-nL})) = O(e^{-nL}).
\]

By decomposing \(L = \ell + 2\pi n\) uniquely for \(\ell \in S^1\) and \(n \in \mathbb{N}\), we arrive at the equivalent system

\[
g(\ell - \varphi + O(e^{-n}), z(\ell + \varphi) + O(e^{-n})) = O(e^{-n}), \quad (\varphi, \ell) \in Q. \tag{5.16}
\]

Lemma 5 shows that \(\Lambda\) and the solution set of (5.16) are diffeomorphic in \(U_\rho(\Lambda_{bif})\) for all \(n\) sufficiently large. Outside this neighbourhood of \(\Lambda_{bif}\), Lemma 7 shows that the solution set of (5.16) is diffeomorphic to \(\Lambda\), since \(\Lambda = Z^{-1}(0)\), zero is a regular value of \(Z\) when restricted to the complement of \(U_\rho(\Lambda_{bif})\), and no additional solutions can appear for (5.16) due to compactness of \(Q\). This completes the proof of Theorem 2.

### 5.3 Isolas of asymmetric 1-pulses

We show that isolas of asymmetric 1-pulses can occur for appropriate functions \(z\).

**Lemma 8** Assume that Hypotheses 6 and 7 are met, and let \(k\) be the number of maxima of the function \(z\). For \(k = 1, 2\), the set \(\Lambda = Z^{-1}(0)\) discussed in Lemma 7 is connected, and therefore does not contain any isolas. For \(k \geq 3\), isolas of asymmetric 1-pulses may exist; an example of a function \(z\) for which an isola exists is given in Figure 8.

**Proof.** Using the procedure outlined above in Figure 7, it is easy to check that isolas cannot exist for functions \(z\) with one or two maxima. We omit a formal proof, which can be obtained using the arguments given in the proof of Lemma 7. Figure 8 gives an example of a function \(z\) that admits an isola. Lemma 7 shows that the drawn isola is necessarily a manifold diffeomorphic to \(S^1\).

### 5.4 \(\mathbb{Z}_2\)-symmetry

Finally, we comment on \(\mathbb{Z}_2\)-equivariant systems for which Hypothesis 5 is met. In this case, Lemma 3 implies that \(g\), and therefore \(z\), are periodic in \(\varphi\) with period \(\pi\). Thus, we can repeat the analysis carried out above for period \(\pi\) instead of \(2\pi\) or, alternatively, simply factor out the \(\mathbb{Z}_2\)-action near the periodic orbits and apply our results to the system on the \(\mathbb{Z}_2\)-orbit space. Either way, we find that the global maxima and minima connect \(\varphi = 0\) to \(\varphi = \pm \pi/2\) and \(\varphi = \pi/2\) to \(\varphi = 0\) and \(\varphi = \pi\).
6 Extensions

We discuss and prove various extensions of our results. Isolas of fronts and reversible systems without a conserved quantity are considered in §6.1 and §6.4, respectively, while higher-dimensional systems and extensions to 2D patterns such as those shown in Figure 2 are treated in §6.2 and §6.3.

6.1 Isolas of fronts

Hypothesis 6 excludes the case illustrated in Figure 5(iii), where the fronts between rolls and the trivial state lie on an isola in parameter space. The expected bifurcation diagram, shown in Figure 6(iii), predicts isolas of symmetric pulses, and such isolas have indeed been observed recently in [3, Figure 24] for the Swift–Hohenberg equation (1.1). We assumed in Hypothesis 6 that the set \( \Gamma \), defined in (2.5) as

\[
\Gamma := \{ (\varphi, \mu) \in S^1 \times J; W^s(0, \mu) \cap W^{uu}(\gamma(\varphi, \mu), \mu) \cap \Sigma_{\text{out}} \neq \emptyset \},
\]

is a graph \((\varphi, z(\varphi))\) and that we can parametrize \(W^s(0, \mu) \cap \Sigma_{\text{out}}\) over \((\varphi, \mu)\). A natural generalization is to assume that there is a smooth function \( G_1 : S^1 \times I \times J \to \mathbb{R} \) so that

\[
(\varphi, v^*) \in W^s(0, \mu) \cap \Sigma_{\text{out}} \iff G_1(\varphi, v^*, \mu) = 0
\]

and \( \Gamma = \{ (\varphi, \mu) \in I \times J; G_1(\varphi, 0, \mu) = 0 \} \), which precludes that \( \Gamma \) intersects \( S^1 \times \partial J \). Note that if Hypothesis 6 is met, we have \( G_1(\varphi, v^*, \mu) = g(\varphi, \mu) - v^* \). The nondegeneracy conditions

\[
\nabla_{(\varphi, \mu)} G_1(\varphi, 0, \mu) \neq 0, \quad \nabla_{(\varphi, v^*)} G_1(\varphi, 0, \mu) \neq 0 \quad \forall (\varphi, \mu) \in \Gamma
\]

(6.1)

guarantee that \( \Gamma \) is a manifold and that \( W^s(0, \mu) \cap \Sigma_{\text{out}} \) is, locally near each element of \( \Gamma \), the graph of a smooth function that depends on \((v^*, \mu)\) or on \((\varphi, \mu)\).

**Theorem 3** In the situation described above, assume that (6.1) and Hypotheses 1-4, and possibly Hypothesis 5, are met. For all points \((\varphi, \mu) \in \Gamma \) near which we can parametrize \( \Gamma \) over \( \varphi \), we also assume that Hypothesis 7 is met. There are then constants \( \rho > 0 \) and \( L_\ast \gg 1 \) so that, for each \( L > L_\ast \) and \( \varphi_0 \in \Theta, (2.1) \) has a symmetric 1-pulse for \( \mu \in J \) that spends time \( 2L \) in the neighborhood \( \mathcal{V} \) near \( \gamma(\cdot, \mu) \) if and only if \( (L, \mu) \in \Gamma_{\varphi_0} \), where \( \Gamma_{\varphi_0} \) is a smooth one-dimensional manifold, which is \( O(e^{-\rho L}) \) close in \( C^1 \) to the manifold

\[
\{(\mu, L) \in J \times (L_\ast, \infty); G_1(L + \varphi_0, 0, \mu) = 0\}.
\]

Furthermore, (2.1) has branches of asymmetric 1-pulses that can be constructed as illustrated in Figure 7 and discussed in §5. If \( \Gamma \) is diffeomorphic to \( S^1 \) and the projection of \( \Gamma \) onto its first component \( \varphi \in S^1 \) is not onto, then \( \Gamma_{\varphi_0} \) is the union of infinitely many disjoint circles, and branches of asymmetric states that emerge and terminate at maxima and minima of \( \Gamma \) connect each of these circles to itself.

In particular, if \( \Gamma \) is an isola, then isolas of symmetric pulses for different \( \varphi_0 \) are not connected, see Figure 6, in contrast to the situation discussed in Theorem 2.

**Proof.** The proofs given in §4 and §5 need only minor modifications which we now outline. For symmetric pulses, the bifurcation equation (4.2) is replaced by

\[
G_1(v^f(L, \varphi_0, \mu), v^s(L, \varphi_0, \mu), \mu) = G_1(\varphi_0 + L + O(e^{-\rho L}), O(e^{-\rho L}), \mu) = 0.
\]

The first nondegeneracy condition in (6.1) allows us to use the implicit function theorem to solve the above equation uniquely near each \( (L, \mu) \) with \( (L \mod 2\pi, \mu) \in \Gamma \). Pitchfork bifurcations from symmetric to asymmetric states occur near elements \((\varphi, \mu) \in \Gamma \) with \( \partial_\varphi G_1(\varphi, 0, \mu) = 0 \): equation (6.1) ensures that \( W^s(0, \mu) \)
can be parametrized via \( v^s = g(\varphi, \mu) \) near such points, and we are therefore back in the setting of §5.1. Branches of asymmetric states can also be continued: (5.10) is now given by
\[
G_1(L + \varphi + O(e^{-\rho L}), O(e^{-\rho L}), \mu) = 0, \quad G_1(L - \varphi + O(e^{-\rho L}), O(e^{-\rho L}), \mu) = 0,
\]
and it is easy to verify that the Jacobian with respect to \((L, \varphi, \mu)\) of the left-hand side has full rank except at pitchfork bifurcation points. This shows that solutions of the asymptotic problem, obtained by omitting the \( O(e^{-\rho L}) \) terms in (6.2), persist for (6.2).

6.2 Higher-dimensional ODEs

We focus now on generalizations of the results in §2 to higher-dimensional systems where \( u \in \mathbb{R}^{2n} \) with \( n > 2 \). In this situation, the parametrization of invariant manifolds is slightly different. We assume that the linearization of the time-2\( \pi \) map of the periodic orbits has two simple real Floquet multipliers at \( e^{\pm 2\pi \alpha(\mu)} \) with \( \alpha(\mu) > 0 \) for all \( \mu \in J \) and that there is an \( \eta > 0 \) so that the remaining Floquet multipliers have modulus either less than \( e^{-2\pi(\alpha(\mu)+\eta)} \) or larger than \( e^{2\pi(\alpha(\mu)+\eta)} \) uniformly in \( \mu \in J \). We can now choose coordinates \((v^c, v^s, v^u, v^{uu})\) near the periodic orbits \( \gamma(x, \mu) \) in the zero energy level set that reflect the uniform spectral decomposition assumed above. In particular, \( W^u(\gamma(x, \mu), \mu) \) is given by \((v^c, 0, 0, v^u, v^{uu})\) in these coordinates, and we can define the section \( \Sigma_{out} \) by requiring \( v^u = \delta \). The key hypothesis is that
\[
W^s(0, \mu) \cap \Sigma_{out} = \{(v^c, v^s, v^u, \delta, v^{uu}) = (v^c, g^s(v^c, v^s, \mu), v^u, \delta, g^{uu}(v^c, v^s, \mu));
(v^c, v^s, \mu) \in S^1 \times B \times J \}
\]
for appropriate smooth functions \( g^s \) and \( g^{uu} \), where \( B \) is a small ball centered at the origin in \( \mathbb{R}^{n-1} \). In particular, zeros of \( g^s \) correspond to heteroclinic orbits between \( u = 0 \) and \( u = \gamma(x, \mu) \). Note that (6.3) requires that the stable manifold \( W^s(0, \mu) \) approaches the unstable manifold \( W^u(\gamma(x, \mu), \mu) \) of the periodic orbits in backward time along the direction associated with the unstable leading Floquet multiplier \( e^{2\pi \alpha(\mu)} \) and that \( W^s(0, \mu) \) intersects \( W^u(\gamma(x, \mu), \mu) \) as transversely as possible.

The results in [37, Theorem 2.1 and (2.8)] show that there is a \( \rho > 0 \) so that Lemma 2 remains true provided (3.2) is replaced by
\[
\begin{align*}
v(-L) &= (\varphi - L + O(e^{-\rho L}), \delta, v^{ss} + O(e^{-\rho L}), O(e^{-\rho L}), O(e^{-\rho L})) \\
v(L) &= (\varphi + L + O(e^{-\rho L}), O(e^{-\rho L}), \delta, v^{uu} + O(e^{-\rho L})) \\
v(0) &= (\varphi, O(e^{-\rho L}), O(e^{-\rho L}), O(e^{-\rho L})).
\end{align*}
\]
To find symmetric 1-pulses, we need to solve
\[
(v^c, g^s(v^c, v^{ss}, \mu), v^{ss}, \delta, g^{uu}(v^c, v^{ss}, \mu)) = (\varphi + L + O(e^{-\rho L}), O(e^{-\rho L}), \delta, v^{uu} + O(e^{-\rho L}))
\]
which, upon solving for \((v^{ss}, v^{uu})\), reduces to
\[
(v^c, g^s(v^c, v^{ss}, \mu)) = (\varphi + L + O(e^{-\rho L}), O(e^{-\rho L})).
\]
A similar reduction works for asymmetric 1-pulses.

We conclude that the results presented in §2 remain true for systems in \( \mathbb{R}^{2n} \), provided we replace all terms of the form \( e^{-2\pi(\mu)\ell} \) by \( O(e^{-\rho \ell}) \). Thus, while we lose information about the specific location of saddle-node and pitchfork bifurcations and about the specific offset of the snaking curves, the results about the bifurcation structure remain valid.
6.3 Snaking in the planar Swift–Hohenberg equation

We consider stationary planar patterns that are periodic and reflection-symmetric in one spatial direction and localized in the other direction, and refer to Figure 2 or [4, Figure 15a] for examples. As shown in Figure 2, snaking is possible for these patterns, and our goal is to show that the results described in §2 apply to such patterns under assumptions similar to those discussed in §6.2.

For the sake of clarity, we outline the framework within which we can study such patterns only for the planar Swift–Hohenberg equation

\[ U_t = -(1 + \Delta)^2 U - \mu U + \nu U^2 - U^3 \quad (6.4) \]

with \((x, y) \in \mathbb{R}^2\). We assume that \(U(x, y)\) is a smooth solution of (6.4) that is periodic in \(y\) with minimal period \(\ell > 0\). We define the wave number \(\kappa = 2\pi/\ell\) and introduce the rescaled variable \(\phi = \kappa y\), which allows us to consider the equation

\[ \partial_t^2 U + 2\kappa^2 \partial_x^2 \partial_y^2 U + \kappa^4 \partial_y^4 U + 2(\partial_x^2 U + \kappa^2 \partial_y^2 U) + (1 + \mu)U - \nu U^2 + U^3 = 0 \]

with \(\phi \in S^1 = [0, 2\pi]/\sim\). For fixed \(\kappa\) and \(\nu\), we write this equation as a first-order system in \(x\) and obtain

\[ u_x = A(\mu)u + N(u) \quad (6.5) \]

where

\[
A(\mu) = \begin{pmatrix}
0 & 1 & 0 & 0 \\
0 & 0 & 0 & 1 \\
0 & 0 & 0 & 1 \\
-\kappa^4 \partial_y^4 - 2\kappa^2 \partial_y^2 - (1 + \mu) & 0 & -2 & -2\kappa^2 \partial_y^2 \\
\end{pmatrix}, \\
N(u) = \begin{pmatrix}
0 \\
0 \\
0 \\
\nu U^2 - U^3 \\
\end{pmatrix}
\]

and

\[ u = (U, U_x, U_{xx}, U_{xxx}, U_{xxxx}) \in Y := H^3_e(S^1) \times H^3_e(S^1) \times H^1_e(S^1) \times L^2_e(S^1) \times L^2_e(S^1) \]

with \(L^2_e(S^1) = \{u \in L^2(S^1) : u(\phi) = u(-\phi)\}\). In other words, for each \(x\), \(u(x)\) is an even periodic function of \(\phi\). We record that \(A(\mu)\) is closed and densely defined on \(Y\) with domain \(Y^1 = H^2_e \times H^3_e \times H^2_e \times H^1_e\), and for each \(\mu > 0\) there is a constant \(C > 0\) so that the resolvent estimate

\[ \|(A(\mu) - ik)^{-1}\|_{L(Y)} \leq \frac{C}{1 + |k|} \quad \forall k \in \mathbb{R} \quad (6.6) \]

holds. Note also that the nonlinearity \(N : Y \to Y\) is smooth due to Sobolev’s embedding theorem. In the remainder of this section, we will denote by \(|\cdot|\) the norm in \(Y\).

Localised structures such as those shown in Figure 2 correspond to homoclinic orbits of (6.5) in \(x\) that connect the trivial state \(u = 0\) to itself in the space \(Y\) of periodic functions in \(\phi\), and we are therefore in a familiar setting. As in §6.2, Lemma 2 is the key technical result that we need to carry forward our analysis, and much of this section will deal with extending Lemma 2 to (6.5).

We first discuss the analogues of Hypotheses 1-7 that we need to analyse symmetric and asymmetric patterns. As shown in [20, §3.3], (6.5) is reversible with reverser \(\mathcal{R}u = \mathcal{R}(u_1, u_2, u_3, u_4) = (u_1, -u_2, u_3, -u_4)\) and conservative with first integral

\[ \mathcal{H}(u, \mu) = \int_0^{2\pi} \left[ u_2 u_4 - \frac{u_2^3}{2} + u_2^2 + \frac{(1 + \mu)u_1^2}{2} - \frac{\nu u_1^3}{3} + \frac{u_1^4}{4} - \kappa^2 (\partial_y u_2)^2 - \kappa^2 (\partial_y u_1)^2 + \frac{\kappa^4 (\partial_y^2 u_1)^2}{2} \right] dy. \]

We infer from (6.6) that the origin is hyperbolic for \(\mu > 0\) as required in Hypothesis 3.
We need to assume the existence of a family $\gamma(x, \mu)$ of symmetric periodic solutions of (6.5) that have zero energy so that $H(\gamma(\cdot, \mu), \mu) = 0$ for all $\mu$ in a closed nontrivial interval $J \subset \mathbb{R}^+$. Furthermore, we assume that the gradient of $H$ does not vanish along the family $\gamma(0, \mu)$ and that the Floquet spectrum of the orbits $\gamma(x, \mu)$ is as described in §6.2, where we remark that [22] ensures that these spectra are well defined for solutions of (6.5). The results proved in [2, 20, 23] show that this analogue of Hypothesis 4 is satisfied by both rolls and hexagons provided $0 < \mu < 1$. If the cubic-quintic nonlinearity $U \mapsto \nu U^3 - U^5$ is used in (6.4), Hypothesis 5 is also satisfied, at least for $0 < \mu < 1$.

It remains to discuss Hypotheses 6-7 which are concerned with the invariant manifolds of the origin $u = 0$ and the periodic orbits $\gamma(x, \mu)$. Even though (6.5) is ill-posed as an initial-value problem, the theory developed in [25, 36] shows that stable and unstable manifolds exist. Moreover, [36, Theorem 5.1] shows that the stable and unstable manifolds of $\gamma$ are foliated by smooth strong stable and strong unstable fibers. Since the zero-energy level set is a smooth manifold near the periodic orbits, we can flatten it. Using the smooth invariant foliations, we can then define coordinates $v = (v^c, v^s, v^{ss}, v^{uu})$ near the periodic orbits $\gamma(x; \mu)$ in the zero-energy level set as in §6.2. In these coordinates, we assume that there exist smooth functions $g^j$ for $j = u, uu, s, ss$ so that we have

$$W^s(0, \mu) \cap \Sigma_{\text{out}} = \{ v = (v^c, g^s, v^{ss}, v^{uu}, v^c, v^{ss})(v^c, v^{ss}, \mu); (v^c, v^{ss}, \mu) \in S^1 \times B \times J \}$$

$$W^u(0, \mu) \cap \Sigma_{\text{in}} = \{ v = (v^c, \delta, g^s(v^c, v^{uu}, \mu), g^u(v^c, v^{uu}, \mu), v^{uu}); (v^c, v^{uu}, \mu) \in S^1 \times B \times J \},$$

where $B$ denotes a small ball in the infinite-dimensional space used to parametrize the strong stable and unstable fibers.

**Lemma 9** Under the above assumptions, there exist positive constants $L_0$ and $\rho$ such that the following is true for all $L > L_0$ and $\varphi \in S^1$: there is a unique solution $v(x, \mu)$ of (6.5), defined for $x \in [-L, L]$, so that

$$v(-L) \in \Sigma_{\text{in}}, \quad v(L) \in \Sigma_{\text{out}}, \quad v(0) = \varphi$$

and $v(x)$ lies in a neighborhood of the periodic orbits $\gamma(x, \mu)$ for all $x \in [-L, L]$. Furthermore,

$$v(-L) = (\varphi - L + O(e^{-\rho L}), \delta, v^{ss} + O(e^{-\rho L}), O(e^{-\rho L}), O(e^{-\rho L}))$$

$$v(L) = (\varphi + L + O(e^{-\rho L}), O(e^{-\rho L}), O(e^{-\rho L}), \delta, v^{uu} + O(e^{-\rho L})) \quad (6.7)$$

$$v(0) = (\varphi, O(e^{-\rho L}), O(e^{-\rho L}), O(e^{-\rho L}), O(e^{-\rho L})).$$

The solution $v(x)$ is smooth in $(\varphi, \mu, L)$, and the error estimates in (6.7) can be differentiated. In addition,

$$v(x, -\varphi) = Rv(-x, \varphi), \quad \varphi \in S^1, \ |x| \leq L$$

and, if Hypothesis 5 holds,

$$v(x, \varphi + \pi) = \kappa v(x, \varphi), \quad \varphi \in S^1, \ |x| \leq L.$$

In particular, the solution $v(x, \varphi)$ is $R$- reversible, with $v(0) \in \text{Fix}(R)$, if and only if $\varphi = 0, \pi$, and $R$- reversible, with $v(0) \in \text{Fix}(\kappa R)$, if and only if $\varphi = \pi/2, 3\pi/2$.

Once this lemma has been proven, analogies of Theorems 1 and 2 can be proven in a manner similar to the finite-dimensional setting, and we omit the details.

**Proof.** We cannot prove the statements within the zero-energy level set as flattening out this manifold destroys the semilinear nature of (6.5) which we need to apply the results on exponential dichotomies in [25]. Instead, we consider (6.5) in the full space $Y$.  

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Hypothesis 4 implies that there is, in fact, a smooth two-parameter family \( \gamma(\cdot, \mu, e) \) of symmetric periodic orbits that are parametrized by \( \mu \) and their energy \( e = H(\gamma(0, \mu, e), \mu) \) for \( e \) near zero. We normalize the family \( \gamma(x, \mu, e) \) of periodic solutions so that \( \gamma(0, \mu, e) \in \text{Fix}(\mathcal{R}) \) for all \((\mu, e)\). For fixed \((\mu, e, \varphi)\), we parametrize the strong stable fiber of \( \gamma(x+\varphi, \mu, e) \) by \( b^+ \) which incorporates the variable \( v^+ \) in the coordinates discussed before Lemma 9 and the time-like direction \( x \). Thus, we denote by \( q^+(x, \mu, e, \varphi, b^+) \) the strong stable fiber of \( \gamma(x+\varphi, \mu, e) \) so that \( q^+ \) is smooth in all arguments and
\[
|\gamma(x+\varphi, \mu, e) - q^+(x, \mu, e, \varphi, b^+)| \leq Ke^{-\rho x}, \quad x \geq 0.
\]
Similarly, \( q^-(x, \mu, e, \varphi, b^-) \) denotes the strong unstable fiber of \( \gamma(x+\varphi, \mu, e) \). From now on, we denote various different constants by the letter \( K \); these different constants will be bounded by a uniform constant that does not depend on any of the parameters or variables involved.

We seek solutions of the form
\[
\begin{align*}
    u^+(x) &= q^+(x, \mu, e^+, \varphi^+, b^+) + w^+(x), \quad x \in [0, L] \\
    u^+(0) &\in \Sigma_{\text{in}} \\
    u^-(x) &= q^-(x, \mu, e^-, \varphi^-, b^-) + w^-(x), \quad x \in [-L, 0] \\
    u^-(0) &\in \Sigma_{\text{out}} \\
    u^+(L) &= u^-(L)
\end{align*}
\]
and refer to Figure 12 for an illustration. Thus, \( u^\pm \) satisfy (6.5) if and only if \( w^\pm \) satisfy
\[
    w^+_x = [A(\mu) + N_u(q^+(x, \mu, e^+, \varphi^+, b^+))] w^++N^+(x, w^+), \quad 0 \leq x \leq L, \tag{6.8}
\]
where \( N^+(x, w^+) := N(q^++w^+)-N(q^+)-N_u(q^+)w^+ = O(|w^+|^2) \), and analogously for \( w^- \). For the Floquet spectrum of the periodic orbits \( \gamma(\cdot) \), we had assumed that there exists a \( \rho > 0 \), independent of \((\mu, e)\), such that the real part of the spectrum, other than the two eigenvalues at zero, is bounded in absolute value from below by \( \rho \). The results in [25] then imply that the linear part of (6.8) has exponential dichotomies \( \Phi^u_+(x, y) \) and \( \Phi^u_+(x, y) \) that are smooth in \((\mu, e, \varphi^+, b^+)\) such that, for \( \epsilon > 0 \) sufficiently small, there exists a \( K > 0 \) with
\[
    \begin{align*}
    |\Phi^u_+(x, y)|_{L(Y)} &\leq Ke^{\epsilon(x-y)}, \quad x \geq y \geq 0 \\
    |\Phi^u_+(x, y)|_{L(Y)} &\leq Ke^{-\rho(y-x)}, \quad y \geq x \geq 0.
\end{align*}
\]
The variation-of-constants formula shows that solutions of (6.8) on \([0, L]\) satisfy the integral equation

\[ w^+(x) = \Phi^u(x, L)a^+ + \int_L^x \Phi^u(x, y)N^+(y, w^+(y)) \, dy + \int_0^x \Phi^c(x, y)N^+(y, w^+(y)) \, dy \]  

(6.9)

and vice versa, where the quantity \(a^+ \in \text{Rg}(\Phi^u(x, L, L))\) is an arbitrary element of the unstable fiber of \(q^+\) at \(x = L\) and represents the final linear displacement off the centre-stable manifold. For each fixed \((\mu, e^+, \varphi^+, a^+, b^+, w^+)\), we wish to view \(F\) as a smooth map from \(C^0([0, L], Y)\), equipped with the weighted norm \(\|w^+\| := \sup_{0 \leq x \leq L} e^{\rho(L-x)}|w^+(x)|\), into itself. Indeed, \(F\) is well defined since

\[ \|F(\mu, e^+, \varphi^+, a^+, b^+, w^+)\| \leq K|a^+| + Ke^{\rho(L-x)} \int_L^x e^{-\rho(y-x)}e^{-2\rho(L-y)}\|w^+\|^2 \, dy + Ke^{\rho(L-x)} \int_0^x e^{(x-y)}e^{-2\rho(L-y)}\|w^+\|^2 \, dy \leq K|a^+| + Ke^{-\rho L} \int_L^x e^{\rho(y)} \, dy + e^{-\rho(L-x)}e^{\epsilon x} \int_0^x e^{(2\rho - \epsilon)y} \, dy \|w^+\|^2 \]

(6.10)

and it is also straightforward to verify that \(F\) is smooth in \((\mu, e^+, \varphi^+, a^+, b^+, w^+)\). In particular, for all sufficiently small \(a^+\), (6.9) has a unique solution \(w^+\), which satisfies \(|w^+(x)| \leq K|a^+|e^{-\rho(L-x)}\) for all \(x \in [0, L]\). We also have

\[ |w^+(L) - a^+| \leq K|a^+|^2, \quad |w^+(0)| \leq K|a^+|e^{-\rho L}. \]

Analogous arguments apply to the function \(w^-\), which satisfies the integral equation

\[ w^-(x) = \Phi^s(x, -L)a^- + \int_{-L}^x \Phi^s(x, y)N^-(y, w^-(y)) \, dy + \int_0^x \Phi^c(x, y)N^-(y, w^-(y)) \, dy \]  

(6.10)

for \(x \in [-L, 0]\), where \(a^- \in \text{Rg}(\Phi^s(x, -L, L))\) is in the stable fiber of \(q^-\) at \(x = -L\). The unique solution of (6.10) satisfies

\[ |w^-(L) - a^-| \leq K|a^-|^2, \quad |w^-(0)| \leq K|a^-|e^{-\rho L}. \]

We now solve the matching condition \(u^+(L) = u^-(L)\) given by

\[ \gamma(L + \varphi^+, \mu, e^+) - \gamma(-L + \varphi^-, \mu, e^-) + a^+ - a^- = O(|a^+|^2 + |a^-|^2 + e^{-\rho L}), \]

(6.11)

where we used the fact that \(|q^+(L, \mu, e^+, \varphi^+, b^-) - \gamma(L + \varphi^+, \mu, e^+)| = O(e^{-\rho L}) \) and analogously for \(q^-\). To solve these equations, we fix \((\varphi, e)\) and set

\[ e^+ = e + \hat{e}, \quad \varphi^+ = \varphi + \hat{\varphi} - L \]

\[ e^- = e - \hat{e}, \quad \varphi^- = \varphi - \hat{\varphi} + L \]

for \((\hat{\varphi}, \hat{e})\) small, so that (6.11) becomes

\[ \gamma(\varphi + \hat{\varphi}, \mu, e + \hat{e}) - \gamma(\varphi - \hat{\varphi}, \mu, e - \hat{e}) + a^+ - a^- = O(|a^+|^2 + |a^-|^2 + e^{-\rho L}). \]

(6.12)

Since the left-hand side vanishes at \((\hat{\varphi}, \hat{e}, a^+, a^-) = 0\), and its derivative with respect to these variables is invertible, we find a locally unique solution \((\hat{\varphi}_0, \hat{e}_0, a^+_0, a^-_0)\) of (6.12), which depends smoothly on the variables \((\mu, e, \varphi, b^-, b^+, L)\). We can also solve the remaining equations

\[ u^+(0) = q^+(0, \mu, e^+, \varphi^+, b^+) + O(e^{-\rho L}a^+) \in \Sigma_{\text{in}} \]

\[ u^-(0) = q^-(0, \mu, e^-, \varphi^-, b^-) + O(e^{-\rho L}a^-) \in \Sigma_{\text{out}} \]
for \((\dot{\varphi}_0, \dot{e}_0, a_{0}^+, a_{0}^-)\) as functions of \((\mu, e, \varphi, b^-, b^+, L)\). We record that \((\dot{\varphi}_0, \dot{e}_0, a_{0}^+, a_{0}^-) = O(e^{-\rho L})\).

We determine now when these solutions are reversible. For given \((\mu, e, \varphi, b^-, b^+, L)\), we define

\[
v^+(x) := R u^-(x), \quad v^-(x) := R u^+(x),
\]

and note that these functions satisfy (6.5) and

\[
v^+(L) = R u^-(L) = R u^+(L) = v^-(L) \\
v^+(0) = R u^- (0) \in R \Sigma_{\text{out}} = \Sigma_{\text{in}} \\
v^-(0) = R u^+ (0) \in R \Sigma_{\text{in}} = \Sigma_{\text{out}},
\]

where we used that \(u^+(L) = u^-(L)\). We want to show that \(v^\pm = u^\pm\) whenever \(\varphi \in \{0, \pi\}\) and \(b^- = R b^+\).

For \(\varphi = 0\) and \(b^- = R b^+\), we have

\[
u^+(x) = q^+(x; \mu, e + \dot{e}_0, \tilde{\varphi}_0 - L, b^+) + \Phi^u(x, L) a^+_0 + \tilde{w}^+(x) \\
u^-(x) = q^-(x; \mu, e - \dot{e}_0, \tilde{\varphi}_0 + L, b^-) + \Phi^s(x, -L) a^-_0 + \tilde{w}^-(x)
\]

for a unique \((\tilde{\varphi}_0, \dot{e}_0, a^+_0, a^-_0)\), where all integrals present in (6.9) are subsumed into \(\tilde{w}^\pm\). Hence,

\[
v^+(x) = R u^- (-x) = q^+(x; \mu, e + \dot{e}_0, -\tilde{\varphi}_0 - L, b^+) + \Phi^u(x, L) R a^+_0 + R \tilde{w}^- (x) \\
v^-(x) = R u^+(x) = q^- (x; \mu, e + \dot{e}_0, \tilde{\varphi}_0 + L, b^-) + \Phi^s (x, -L) R a^-_0 + R \tilde{w}^+ (x).
\]

In the above, we have used \(\mathcal{R} q^+ (x; \mu, e + \dot{e}_0, \tilde{\varphi}_0 - L, b^+) = q^- (x; \mu, e + \dot{e}_0, -\tilde{\varphi}_0 + L, R b^+)\) together with \(b^- = R b^+\), and similarly for \(q^\pm\). We have shown above that \((\tilde{\varphi}_0, \dot{e}_0, a^+_0, a^-_0)\), and therefore the functions \(\tilde{w}^\pm\), are uniquely determined by \((\mu, \varphi, e, b^+, b^-, L)\). On the other hand, a comparison of the above expressions for \(u^\pm\) and \(v^\pm\) shows that both \((\tilde{\varphi}_0, \dot{e}_0, a^+_0, a^-_0)\) and \((-\tilde{\varphi}_0, -\dot{e}_0, R a^+_0, R a^-_0)\) give solutions for the same data \((\mu, \varphi, e, b^+, b^-, L)\) which means that \(\dot{e}_0 = \tilde{\varphi}_0 = 0\) and \(u^\pm = v^\pm\) as claimed. This establishes that \(u^\pm\) are reversible whenever \(\varphi = 0\) and \(b^- = R b^+\), and a similar argument gives reversibility when \(\varphi = \pi\).

Finally, we need to show that we can choose \(e\) so that the solutions \(u^\pm\) lie in the zero energy set. Using that \(\mathcal{H}\) is conserved along solutions, we find

\[
\mathcal{H}(u^+(0)) = \mathcal{H}(u^+(L)) = \mathcal{H}(q^+(L, \mu, e, \varphi - L, b^+) + O(e^{-\rho L})) \\
= \mathcal{H}(\gamma(\varphi, \mu, e) + O(e^{-\rho L})) = \mathcal{H}(0) + O(e^{-\rho L})
\]

and therefore \(\frac{d}{de} \mathcal{H}(u^+(0)) \neq 0\) near \(e = 0\) by hypothesis. Thus, we can solve near \(e = 0\) for the unique solution \(e = e(\mu, \varphi, b^+, L)\) that will guarantee that \(\mathcal{H}(u^+) = 0\).

We are now in a position to transfer the results obtained in \(Y\) to the flattened zero-energy level set by applying out the coordinate transformations described briefly before Lemma 9, and we omit the details.

### 6.4 Reversible systems without first integral

We assumed in Hypothesis 2 that the underlying ODE is conservative. For symmetric pulses, this assumption is not needed, and Theorem 1 is true for systems that satisfy only Hypotheses 1, 3-4, and 6-7. More precisely, we assume that there is a smooth two-parameter family \(\gamma(\cdot, \mu, e)\) of symmetric periodic orbits that are parametrized by \(\mu\) and a second internal parameter \(e\). To find homoclinic orbits, we can utilize the Shilnikov variables that were discussed in the proof of Lemma 9 in §6.3 to characterize solutions that pass near the two-parameter family of symmetric periodic orbits. Symmetric orbits can then be found by following the
proofs in §4 except that we now solve simultaneously not only for \((L, \mu)\) but also for the variable \(e\) that encodes which of the reversible orbits \(\gamma(\cdot, \mu, e)\) the homoclinic orbit follows.

To construct asymmetric pulses, we note that a first integral in \(\mathbb{R}^4\), say, allows us to reduce each of the three-dimensional sections \(\Sigma\text{in}\) and \(\Sigma\text{out}\) to a two-dimensional section. If we do not have a first integral, we need two additional free variables to be able to solve in each section in the direction corresponding to \(e\) along the one-parameter family \(\gamma(x, \mu, e)\) of reversible periodic orbits. In one section, we can simply use the parameter \(e\) itself. For non-conservative systems, no other free variables are available, and stationary asymmetric pulses should therefore not exist. Instead, asymmetric pulses should begin to move, and we are led to consider the Swift–Hohenberg equation in a moving frame \(y = x - ct\),

\[
U_t = -(1 + \partial_y^2)U + c\partial_y U - \mu U + \nu U^2 - U^3,
\]

and seek stationary localised states \(U(y)\) for this equation, where \(c\) provides now an additional free parameter. Using the unique periodic solution of the adjoint variational equation around each reversible periodic orbit, we find that the equation for \(e\) in Fenichel’s normal form becomes \(e_y = c + \text{h.o.t.}\). Thus, we can use \(e\) and \(c\) to solve in the remaining directions in each section, and we conclude that asymmetric pulses exist exponentially close to the sets described by Lemma 7 for wave speeds that are \(O(e^{-\eta L})\).

7 Discussion

In this paper, we identified conditions on the global bifurcation structure of heteroclinic orbits between equilibria and periodic orbits that guarantee that we can glue connecting orbits together to construct both symmetric and asymmetric orbits that are homoclinic to the underlying equilibrium. We also showed that snaking occurs provided the global structure of fronts is as shown in Figure 5(i)-(ii), while the structure sketched in Figure 5(iii) leads to infinitely many isolas of pulses.

Our results can be used to predict all branches of asymmetric pulses, including isolas, once the snaking curves of the symmetric states have been computed. Indeed, Figure 7 illustrates how these branches can be found from the function \(z\), while (1.4) establishes that the branch \(\mu_*(L, \varphi_0)\) of symmetric pulses is exponentially close to the graph \(\mu = z(L + \varphi_0)\) of \(z\), so that we can reliably determine \(z\) through numerical computations of the branch \(\mu_*(L, \varphi_0)\) of symmetric pulses.

**PDE stability.** We now comment on temporal stability of symmetric and asymmetric pulses in the context of self-adjoint partial differential equations (PDEs) such as the Swift–Hohenberg equation (1.1). Numerical computations of the first few eigenvalues of symmetric pulses of (1.1) have been published in [3, Figures 10-11], and our goal here is to compare these computations with theoretical predictions.

The spectrum of each front is the union of isolated eigenvalues and the essential spectra \(\Sigma_{\text{ess}}\) and \(\Sigma_{\text{rolls}}\) of the trivial state and the asymptotic rolls, respectively. We assume that the front spectrum looks as shown in the top panel of Figure 13(i). In particular, we assume that \(\Sigma_{\text{ess}}\) lies on the negative real axis \(\mathbb{R}^-\), at some distance away from zero, and that \(\Sigma_{\text{rolls}}\), which emerges from \(\lambda = 0\), lies in \(\mathbb{R}^-_0\). The origin \(\lambda = 0\) is a branch point of the dispersion relation of the rolls. It remains to discuss the point spectrum of the fronts. Inspecting a typical bifurcation diagram such as Figure 5(i), we may expect that there is a single eigenvalue associated with the fronts that moves forth and back through zero at each fold bifurcation: note, however, that once the “eigenvalue” lies on the negative real axis, it is contained in \(\Sigma_{\text{rolls}}\), and there may no longer be a genuine eigenfunction associated with it. The Evans function\(^4\) provides a mechanism for identifying the saddle-node

\(^{4}\text{We refer to [33] for a survey}\)
Figure 13: Panel (i) shows the anticipated PDE spectra of fronts [top] and symmetric 1-pulses [bottom]: \( \Sigma_{\text{ess}} \) denotes the essential spectrum of the trivial state \( U = 0 \); the essential spectrum of the asymptotic rolls for fronts breaks up into \( O(L) \) eigenvalues for pulses; if the front has a single saddle-node eigenvalue, then we expect that pulses have two eigenvalues that move forth and back as we move along the snaking curve and cause saddle-node and pitchfork bifurcations. The expected movement of the saddle-node [solid circles] and pitchfork [open circles] eigenvalues is shown in more detail in panels (ii) and (iii), corresponding to \( g \mu > 0 \) and \( g \mu < 0 \), along branches of symmetric [solid] and asymmetric [dotted] 1-pulses, where we assumed that symmetric pulses are stable along the lower branch.

eigenvalue even when it lies in \( \mathbb{R}^- \). The results in [15, 16] show that the Evans function \( E(\gamma, \mu) \) is analytic in \( \gamma = \sqrt{\lambda} \) near the origin\(^5\) and that roots \( \gamma \) in \( \operatorname{Re} \gamma > 0 \) of \( E(\gamma, \mu) \) correspond to genuine eigenvalues in \( \mathbb{R}^+ \). Thus, saddle-node eigenvalues can be tracked as zeros of \( E \) even for \( \lambda < 0 \). We also expect that \( \gamma_0 = 0 \) is a zero of the Evans function for all \( \mu \) due to translation invariance and that this root is simple except at folds.

Next, we consider symmetric pulses along the snaking curve which are constructed by gluing together a front \( U_f(x) \) and its symmetric counterpart \( U_f(-x) \). We argue that the spectra of symmetric pulses should look as shown in the bottom panel of Figure 13(i). The essential spectrum of the linearization about a symmetric pulse is given by the spectrum \( \Sigma_{\text{ess}} \) of \( U = 0 \). The essential spectrum \( \Sigma_{\text{rolls}} \) of rolls is not in the spectrum of the pulses but instead breaks up into \( O(L) \) eigenvalues whose union converges to \( \Sigma_{\text{rolls}} \) in the symmetric Hausdorff distance as \( L \to \infty \) [28, 34, 35]. Preliminary expansions of the Evans function associated with the pulses indicate that the root \( \gamma_0 = 0 \) of the front persists as a simple eigenvalue of the pulses, which is in agreement with the numerical computations in [3, Figure 10].

For parameter values \( \mu \) for which the front has an unstable saddle-node eigenvalue \( \lambda_f > 0 \), [34] shows that symmetric pulses will have two eigenvalues, belonging to even and odd eigenfunctions, respectively, within \( O(e^{-\eta L}) \) distance of \( \lambda_f \). At saddle-node and pitchfork bifurcations, these eigenvalues will cross through the origin into the negative real axis, where they collide with the \( O(L) \) absolute eigenvalues which emerge from the absolute spectrum of the rolls and are situated at approximately \( O(1/L) \)-distance from each other. The spectral picture shown in [3, Figure 10] suggests that the saddle-node and pitchfork eigenvalues can be distinguished from these absolute eigenvalues. Our numerical computations, shown in Figure 14, indicate that this is not the case: instead, the saddle-node eigenvalue assumes the role of the first absolute eigenvalue.

We do not know whether this is an artifact of the discretization or a genuine effect of the PDE operator but believe it is the latter.

Finally, we comment on the expected spectra of the asymmetric pulses that bifurcate at pitchfork bifurcations from the symmetric pulses. As illustrated in Figure 13(ii)-(iii), we expect asymmetric pulses to be unstable near onset, and we shall assume that symmetric pulses are stable along the lower branch. In the case of Figures 5(i) and 6(i), the asymmetric pulses do not undergo additional fold bifurcations and should therefore

\(^5\)To keep with standard notation, \( \gamma \) stands for \( \sqrt{\lambda} \), and not for the periodic orbits, in this discussion.
Figure 14: Panel (i) shows the numerically computed saddle-node eigenvalue $\lambda_{\text{sn}}$ of the symmetric pulses of (1.1) for $\nu = 1.6$ along the snaking curve. Panel (ii) contains a schematic plot of the numerically computed saddle-node eigenvalue $\lambda_{\text{sn}}$ from panel (i) [upper-left dotted] and an absolute eigenvalue $\lambda_{\text{abs}}$ [lower-right dotted], plotted together with the eigenvalue $\lambda_f$ of the front [solid blue] and the theoretical absolute eigenvalues in $\Sigma_{\text{abs}}$ [horizontal lines]. The computations in (i) suggest that the saddle-node eigenvalue becomes the first absolute eigenvalue once it reaches $\Re \lambda < 0$.

be unstable across the branch. In contrast, in the situation shown in Figures 5(ii) and 6(ii), the asymmetric pulses undergo additional folds as indicated in Figure 13(ii)-(iii), and the two critical eigenvalues are expected to be either both unstable or both stable along the middle branch.

**Formal front interaction equations.** To gain insight into the stability properties of asymmetric pulses, we discuss a formal approach in the spirit of collective coordinates, which attempts to capture the temporal interaction of the fronts that make up the pulses we constructed. We emphasize that this approach is, in the current context, entirely formal\(^6\). Starting point is the assumption that the underlying PDE

$$U_t = -(1 + \partial_x^2)^2 U + F(U, \mu)$$

has an invariant manifold consisting of possibly time-dependent pulse-like solutions of the type shown in the bottom panel of Figure 13(i). We parametrize this manifold by $U = U(L, \varphi)$, where $2L$ is the spatial width of the roll plateau and $\varphi$ is the phase of the solution, relative to the rolls, at $x = 0$. Substitution into the PDE gives

$$U_{LL} + U_{\varphi\varphi} + (1 + \partial_x^2)^2 U + F(U, \mu) \tag{7.1}$$

where $(L, \varphi)$ depends only on $t$. In §3-5, we showed that stationary solutions of (7.1) with $(L_t, \varphi_t) = 0$ are determined by the matching conditions (5.2). If $(L_t, \varphi_t)$ is small, then the left-hand side of (7.1) is, for each fixed $t$, a small inhomogeneous perturbation of the ODE on the right-hand side, and we arrive at a matching condition similar to (5.2) which then uniquely determines $(L_t, \varphi_t)$; this is one way of deriving the pulse-interaction equations in [8, 40] and [33, §5.3]. In our situation, $(L_t, \varphi_t)$ is not small, and the resulting matching condition could be complicated. Without any rigorous justification, we assume that the contribution of the left-hand side to (5.2) is linear. Assuming furthermore, for simplicity, that $g(\varphi, \mu) = z(\varphi) - \mu$, then a comparison of the parametrization by $(L, \varphi)$ with (3.2) leads to the dynamical system

\[
L_t + \varphi_t = a[z(L + \varphi) - \mu] + O(e^{-\rho L}) \tag{7.2} \\
L_t - \varphi_t = a[z(L - \varphi) - \mu] + O(e^{-\rho L})
\]

for the temporal evolution of $(L, \varphi)$. Adding and subtracting the two equations in (7.2), we obtain

\[
L_t = \frac{a}{2}[z(L + \varphi) + z(L - \varphi) - 2\mu] + O(e^{-\rho L}) \tag{7.3} \\
\varphi_t = \frac{a}{2}[z(L - \varphi) - z(L + \varphi)] + O(e^{-\rho L}).
\]

\(^6\)Situations in which interaction equations can be justified are discussed in [8, 40] and the references therein.
In the Swift–Hohenberg equation, the trivial state invades rolls for all sufficiently large \( \mu \), while rolls invade the background state for \( \mu \) sufficiently small. This requires \( a > 0 \), so that \( L_t < 0 \), and we therefore focus on the equation

\[
\begin{align*}
L_t &= z(L + \varphi) + z(L - \varphi) - 2\mu + O(e^{-\rho L}) \\
\varphi_t &= z(L - \varphi) - z(L - \varphi) + O(e^{-\rho L}),
\end{align*}
\]

obtained by rescaling time in (7.3) for \( a > 0 \). The linearization of (7.4) about an equilibrium has eigenvalues \( \lambda_{\pm} = z'(L \pm \varphi) + O(e^{-\rho L}) \). In particular, we can read off the PDE eigenvalues predicted by this formal approach from the geometric construction of asymmetric pulses that we outlined in Figure 7. In the situation of Figure 13(ii)-(iii), these heuristic arguments predict that the asymmetric pulses are stable on the middle branch, since we have \( z'(L \pm \varphi) < 0 \) along the middle branch, while they have one unstable eigenvalue on the other two branches in agreement with Figure 13(ii)-(iii).

The preceding formal analysis also predicts that symmetric pulses with \( \varphi_0 \in \{0, \pi\} \) have two eigenvalues exponentially close to \( z'(L + \varphi_0) \), which agrees with the discussion above that led to Figure 13(i). The additional correlation between the derivative of the snaking curve and the fold and pitchfork eigenvalues of symmetric pulses is, at least qualitatively, reflected in the numerical computations of [3, Figure 10].

While the arguments outlined above are formal, we expect that our predictions for the location of eigenvalues of asymmetric pulses can be proved, or disproved, rigorously near each fold on the asymmetric bifurcation curves as we can utilize centre-manifold reductions, but we have not yet pursued this further.

**Swift–Hohenberg equation.** In §4, we argued that the Hypotheses 1-3 we imposed in §2 are satisfied for the Swift–Hohenberg equations (1.1) and (1.2), and the latter equation also satisfies Hypothesis 5. Near onset to rolls, the results in [18] indicate that the remaining Hypotheses 4 and 6-7 are also met, and our results show that snaking occurs as long as these assumptions hold. Since the reverser is generated by \( x \mapsto -x \), (1.1) should therefore exhibit two branches of even localised states corresponding to \( \varphi_0 = 0, \pi \): solutions on the first branch have a maximum at \( x = 0 \), while solutions on the second branch have a minimum at \( x = 0 \). Similarly, the cubic-quintic Swift–Hohenberg equation (1.2) will exhibit two branches of symmetric pulses: one branch corresponds to even solutions \( U(x) \) and their symmetric counterparts \( -U(x) \), while the other branch consists of odd solutions and their symmetric counterparts.

In conclusion, the theory presented here agrees with the earlier numerical results for both symmetric and asymmetric pulses that are summarised in Figure 1. Isolas of symmetric pulses were found in [3, Figure 24], which is in agreement with Theorem 3. Structures that are localised in one direction and periodic in the other direction were found recently in [1] for the planar Swift–Hohenberg equation. The numerically computed snaking curve resembles Figure 5(ii), and Theorem 2 predicts several interesting branches of asymmetric states, see Figure 6(ii). The conjecture stated above asserts that the bifurcating asymmetric branches are stable along the middle branch. Subsequent to our analysis, we have meanwhile computed the predicted asymmetric branches [1], and initial computations indicate that these solutions are indeed stable on the middle branch.

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