

Determination of the single-ion anisotropy energy in a $S = 5/2$ kagome antiferromagnet using x-ray absorption spectroscopy.

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We report x-ray absorption and x-ray linear dichroism measurements at the Fe $L_{2,3}$ edges of the geometrically frustrated systems of potassium and hydronium iron jarosite. Comparison with simulated spectra, involving ligand-field multiplet calculations modelling the $3d-2p$ hybridization between the iron ion and the oxygen ligands, has yielded accurate estimates for the ligand metal-ion hybridization and the resulting single-ion crystal field anisotropy energy. Using this method we provide an experimentally verified scenario for the appearance of a single-ion anisotropy in this nominally high-spin $3d^5$ orbital singlet 6S system, which accounts for features of the spin-wave dispersion in the long-range ordered ground state of potassium iron jarosite.

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I. INTRODUCTION

The jarosite group of minerals has received much attention for being close to ideal realizations of the kagome antiferromagnet. The kagome topology, which is named after a Japanese basket weaving pattern, frustrates the antiferromagnetic Néel ordering¹. Due to the resulting “underconstraint” neither in the quantum case of $S = 1/2$ nor in the classical limit is the kagome antiferromagnet expected to show a symmetry-breaking transition, even at $T = 0$ K^{2,3,4}. Physical realizations of this system are highly valued because they allow for the investigation of the Mott insulating phase in the absence of Néel-like magnetic order. The jarosite group of general stoichiometry $AM_3(\text{SO}_4)_2(\text{OH})_6$, where A can be—amongst others—K, Na, Rb, Ag, NH_3 , or H_3O , provides physical models for the kagome antiferromagnet with spin $S = 3/2$ and $5/2$, where the M site is occupied by trivalent $\text{Cr}^{5,6,7}$ and $\text{Fe}^{5,8,9,10,11,12,13,43}$, respectively. $M = \text{V}$ is a $S = 1$ system with ferromagnetic near-neighbor interactions^{14,15}. Fig. 1 illustrates the jarosite structure of well separated kagome layers consisting of MO_6 octahedra. Potassium iron jarosite is representative for most of the iron analogues. In this system a transition to a non-collinear long-range ordered state is observed at 64 K, despite that the magnitude of the Curie-Weiss temperature, Θ_{CW} , is as high as -800 K^{11,16}. The ground state has shown to be a so-called $q = 0$ state of positive chirality, in which the spins lie within the kagome plane and all point either in or out of each shared kagome triangle^{9,12,13,17,18}. The large figure $\Theta_{\text{CW}}/T_{\text{N}} = 12.5$ is an indication of relatively strong geometric frustration¹⁹, but clearly the observed long-range ordered ground state points to additional terms in the Hamiltonian of this kagome antifer-

romagnet, beyond near-neighbor exchange. It has been shown that the $q = 0$ ground state can arise due to magnetic anisotropies such as easy-plane single-ion crystal field (CF) anisotropy $D_z \hat{S}_z^2 - E_{xy}(\hat{S}_x^2 - \hat{S}_y^2)$,^{20,21} where D_z is the zero-field splitting parameter, or a Dzyaloshinsky-Moriya interaction (DMI)^{21,22} $D_{ij} \hat{S}_i \times \hat{S}_j$. The DMI term has shown to be symmetry allowed on the kagome lattice and has been argued to prevail over the CF anisotropy in the spherical ($L = 0$) $3d^5$ Fe^{3+} ion²². Fits to the spin-wave dispersion curves measured with neutron spectroscopy on single crystal²³ and powder samples²⁴ of potassium jarosite have found to be slightly better in a model with DMI compared to the model with CF anisotropy. The evidence so far has, however, been inconclusive.

Given the similarity in crystal structure of potassium (iron) jarosite and hydronium (iron) jarosite the above effects could be expected to be equally important in the latter. However, hydronium jarosite has a spin-glass ground state with a freezing temperature of $T_g = 17(2)$ K²⁵. The Weiss temperature for hydronium jarosite is less accurately known (-700 to -1400 K) because the inverse susceptibility does not enter a linear regime sufficiently far below the decomposition temperature of the compound¹⁰. It has been suggested that the spin-glass state is the result of proton transfer from the hydronium groups to the $\text{Fe}-(\text{O}_{\text{eq}}\text{H})-\text{Fe}$ super-exchange mediating hydroxy groups¹⁶. One indication that this could happen might be that the OH groups at the crystallographically analogous location in the related kagome compound zinc paratacamite can form a muonium state by chemically binding a positive muon²⁶. The aim of the work described here is to investigate the origins of the magnetic anisotropy (of either single-ion or DMI character) in the

iron jarosites, and to understand the essential differences between the salts such as potassium jarosite with a long range ordered ground state and hydronium jarosite.

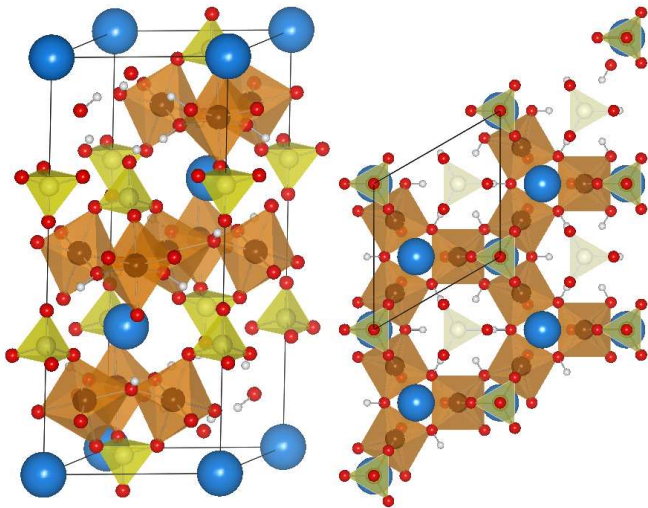


FIG. 1: (color online) The jarosite structure viewed perpendicular to the c axis (left picture) and along the c axis (right picture), revealing the kagome network of magnetic ions. The MO_6 octahedra are drawn in transparent brown with the M^{3+} ions at the center in dark brown. The SO_4 tetrahedra are transparent yellow. The large spheres at the origin of the unit cells in blue are the A site ions, and the hydrogen of the OH groups are drawn as small white spheres.

Using Mössbauer spectroscopy^{27,28} the spin state of the Fe^{3+} in the iron jarosites was found to be $S = 5/2$, thus nominally high-spin $3d^5$. The free-ion $3d^5$ state is 6S , i.e. $L = 0$. Hence, there is no orbital angular momentum due to the spin-orbit coupling in the presence of a crystal field as long as the system is high-spin $3d^5$. This holds at least up to third order perturbation theory of the trigonal crystal-field distortions, spin-spin and spin-orbit couplings⁵⁶. Even in a D_{4h} crystal field the $3d^5$ $S = 5/2$ spin should therefore be practically isotropic apart from the possible effects of dipole-dipole interactions. Both DMI and single-ion anisotropy arise only when the spin-orbit operator $L \cdot S$ can induce a finite orbital angular momentum^{29,30}. That $\langle L \rangle > 0$ in the iron jarosites is clear from Curie-Weiss fits to the high temperature magnetic susceptibility which indicate μ_{eff} between 6.2 and 6.7 μ_B in potassium iron jarosite¹³ and 6.6(2) μ_B in the hydronium analogue^{10,31} (compared to a spin only value of the effective moment of 5.92 μ_B). Clearly, this opens up the possibility for both DMI and single ion anisotropies. However, the origin of the orbital angular momentum itself has so far received little attention.

A very strong crystal field, such as in iron phthalocyanine³² can give rise to a low-spin ground state in which the spin-orbit coupling will in general reinstate part of the orbital angular momentum of the free-ion configura-

tion which is, in zero order, quenched by the crystal field. It could be that in weaker crystal fields, such as those present in iron oxides, some orbital angular momentum appears due to a mixing in of the low-spin state into the ground state. The crystal field alone does, however, not provide matrix elements between high-spin and low-spin states. The most likely explanation is therefore that an Fe^{2+} configuration mixes in due to charge transfer³³ from the surrounding ligands. For hematite, e.g., the effective valence of the Fe ion has been estimated at 2.43 due to charge transfer from the oxygen ligands³⁴. These effects have earlier been shown to explain zero field splitting in Mn^{2+} and Fe^{3+} substituted crystals^{35,36} and in the Mn^{2+} compounds MnF_3 ³⁷, $MnPS_3$ and $MnPSe_3$ ³⁸, despite the lower covalence of the fluoride and phosphorus-containing ligands compared to oxygen.

Transition metal $2p \rightarrow 3d$ x-ray spectroscopy studies have played a key role in the understanding of the effect of charge-transfer and crystal-field in determining the electronic and magnetic properties in transition metals and their compounds^{33,39,40,41}. We have measured the x-ray absorption and x-ray linear dichroism (XLD) at the Fe $L_{2,3}$ ($2p \rightarrow 3d$) edges on single-crystalline and powder samples of potassium and hydronium iron jarosite. By comparison with atomic multiplet calculations, taking into account ligand-field effects and multiple ionic configurations, we obtain accurate values for the charge transfer and single-ion anisotropy in the $S = 5/2$ kagome antiferromagnet. The results are compared to magnetic susceptibility and neutron spectroscopy measurements on in particular potassium iron jarosite.

II. EXPERIMENTAL

Single crystals of potassium iron jarosite, $KFe_3(OH)_6(SO_4)_2$, were grown using a hydrothermal reduction-oxidation method as described in Ref. 16; 4.88 g (28.0 mmol) K_2SO_4 and 2.2 mL (40 mmol) H_2SO_4 were dissolved in 50 mL distilled water, and transferred into a 125 mL PTFE liner of a stainless steel bomb. 0.56 g (10 mmol) iron wire with a diameter of 2 mm was added to the solution. The bomb was placed in an oven at 202°C for 4 days, then cooled down to room temperature at a rate of 0.3°C/min. The precipitate was washed, filtered and dried, yielding 0.37 g, which is 22% based on Fe. In order to obtain sufficiently large crystals it was important that the surface area of the solution exposed to air inside the vessel was minimized. This was achieved using a glass container which fitted inside the PTFE liner, with only a small hole at the top. Using this method single crystals of up to 0.7 mm in diameter were obtained, which was sufficiently large for the application in polarized x-ray spectroscopy. A number of crystals were characterized using single-crystal x-ray diffraction. These crystals were identified as potassium iron jarosite and face-indexed.

The original solvothermal synthesis method^{42,43}

for the jarosites was used for the preparation of $\text{H}_3\text{OFe}_3(\text{OH})_6(\text{SO}_4)_2$. In this preparation 6.6 g (22 mmol) of $\text{Fe}_2(\text{SO}_4)_3 \cdot 5\text{H}_2\text{O}$ was dissolved in 50 mL water. This solution was transferred to a 125 mL PTFE liner of a stainless steel bomb. The solution was heated in the bomb to 140°C for 12 hours. The hydronium iron jarosite which precipitated during the reaction was washed, filtered and dried, yielding ~ 0.27 g of product.

Polarized soft x-ray absorption measurements were carried out using the liquid helium cryostat, housed in the high-field superconducting magnet on ID08 at the European Synchrotron Radiation Facility (ESRF) in Grenoble, France. Total-electron-yield spectra were measured by recording the drain current from the sample as a function of photon energy in the region of the Fe $L_{2,3}$ absorption edges. Potassium iron jarosite is strongly insulating, so the resonance enhanced total-electron-yield signal at the iron edge is much smaller than the signal from the surrounding metal of the sample holder. To avoid this source of background intensity a potassium iron jarosite single crystal of 0.7 mm diameter was attached using silver glue to the end of a graphite tip, which in turn was glued to the end of a standard ESRF sample holder. The sample was aligned in the x-ray beam with the crystallographic c axis pointing in the horizontal plane and perpendicular to the incident beam. This allowed us to measure absorption spectra with the incident x-ray polarization both perpendicular and parallel to the c axis. Spectra were recorded with alternating horizontal and vertical polarization, and for each polarization averages were taken over at least eight scans. In this way the polarization dependent spectra were taken at a number of temperatures between 20 and 290 K, in zero external field.

Isotropic absorption spectra at the Fe $L_{2,3}$ edges were obtained on beamline 5U.1 of the Synchrotron Radiation Source (SRS) at Daresbury Laboratory, UK, using potassium and hydronium iron jarosite powder dispersed over double-sided UHV-compatible carbon tape.

III. RESULTS

Figure 2 shows the spectra obtained from single-crystal and powder samples of potassium iron jarosite at 290 K. The upper curves in Fig. 2 display the x-ray absorption spectra from a single crystal measured on beamline ID08 at the ESRF, with the incident x-ray polarization parallel to the c axis, I_c (black circles), and perpendicular to the c axis, I_{ab} (red crosses). The curves in the middle of the figure give the corresponding isotropic absorption spectrum $I_{\text{iso}} = \frac{2}{3}I_{ab} + \frac{1}{3}I_c$ (black line) and the isotropic spectrum directly measured from a powder sample on station 5U.1 at the SRS (broken red line). The curve at the bottom gives the x-ray linear dichroism I_{xld} defined as $I_{ab} - I_c$ (thin black line). The $2p \rightarrow 3d$ transition consists of two edges, the first one at ~ 709 eV (L_3 edge) corresponds to the creation of a $2p_{3/2}$ core hole. The

transition giving rise to a $2p_{1/2}$ core hole (L_2 edge) is at ~ 12 eV higher photon energy due to the $2p$ spin-orbit coupling. The inset of Fig. 2 shows the I_c and I_{ab} spectra in close up at the L_3 edge, which is the most significant region of the spectrum.

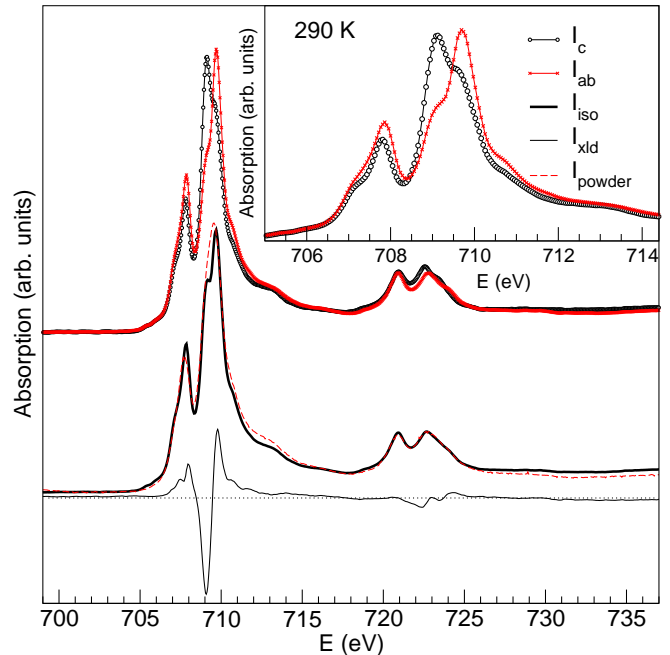


FIG. 2: (color online) x-ray absorption spectra of potassium iron jarosite at 290 K measured with x-ray polarization parallel and perpendicular to the c axis, I_c (black circles) and I_{ab} (red crosses), respectively. The lower trace shows the corresponding values of the isotropic absorption, $I_{\text{iso}} = \frac{2}{3}I_{ab} + \frac{1}{3}I_c$ (black line) and the x-ray linear dichroism, $I_{\text{xld}} = I_{ab} - I_c$ (thin black line). Also shown I_{iso} from a powder sample measured on 5U.1 at SRS Daresbury with a lower energy resolution (broken, red line). The inset shows a close up of I_c and I_{ab} at the L_3 edge.

As seen in Fig. 3, the structure at the L_3 edge becomes more prominent in the isotropic spectrum as the temperature is lowered. At each stabilized temperature, the reproducibility of up to eight spectra demonstrated the absence of any temporal variation in the spectral line shape or intensity, therefore charging effects and drift in sample alignment can be eliminated as the cause of the observed temperature dependence. These isotropic spectra obtained by weighted averaging of I_{ab} and I_c are in good agreement with the spectra from powder samples measured for the SRS at all temperatures. That this change in line shape occurs in the powder averaged spectrum might be an indication of a slight change in the electronic structure as the temperature is lowered. Only very small changes were observed in the XLD spectral line shape on cooling the sample, as shown in Fig. 4. This change in the spectral line shape is accompanied by

a slight reduction in the integrated XLD spectrum⁴⁴

$$\int I_{\text{xld}}(\omega)d\omega / \int I_{\text{iso}}(\omega)d\omega \quad (1)$$

corresponding to a small change of the total quadrupole moment of 0.05(2). This might be interpreted as that the Fe 3d shell gains between 3 and 7% net $z^2 - r^2$ character at base temperature.

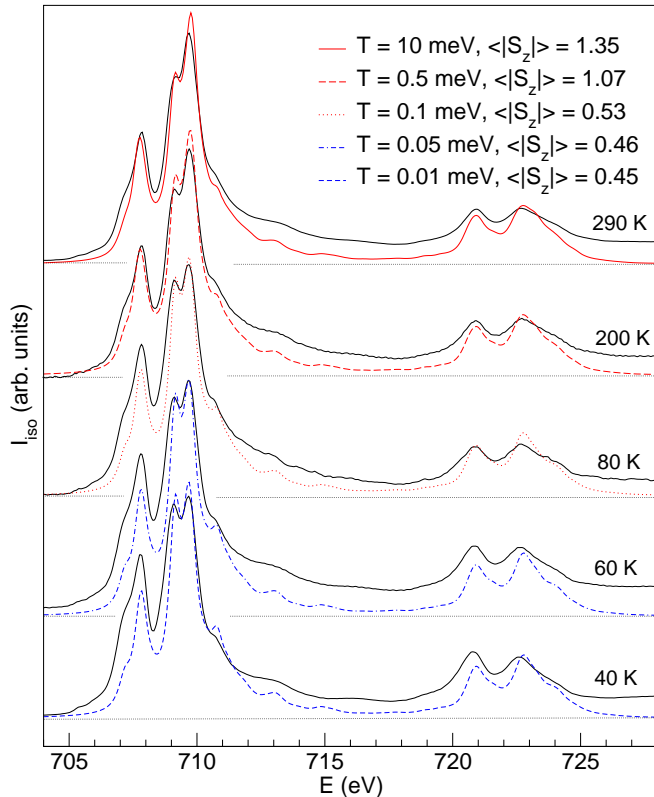


FIG. 3: (color online) Temperature dependence of the isotropic x-ray absorption measured at the Fe $L_{2,3}$ edges from a synthesized potassium iron jarosite single crystal (black lines), compared to calculated spectra (as discussed in Secs. IV and V) with simulated temperatures as listed in the legend.

The room temperature spectra of hydronium iron jarosite were identical to the spectra obtained on potassium iron jarosite within the experimental resolution as measured at the SRS (the red, broken line in the main panel of Fig. 2 for comparison).

IV. LIGAND-FIELD MULTIPLET CALCULATIONS

Due to the strong electrostatic interactions between the $2p$ core hole and the $3d$ levels in the final state, the $2p \rightarrow 3d$ absorption spectrum is not simply proportional to the density of unoccupied $3d$ levels⁴⁵ as a function of energy. For ferromagnets and magnetically soft materials, where all the magnetic moments can be aligned by

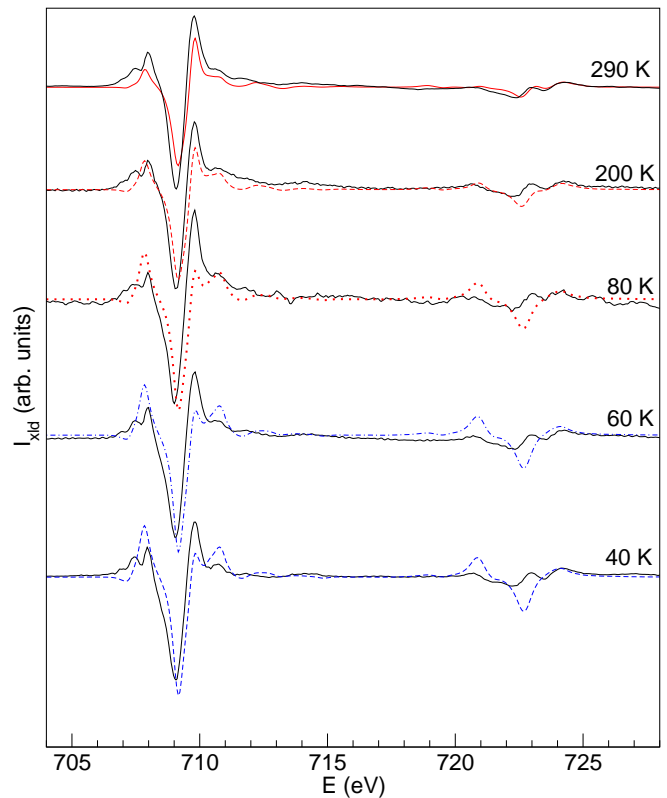


FIG. 4: (color online) Temperature dependence of the x-ray linear dichroism measured at the Fe $L_{2,3}$ edges from a synthesized potassium iron jarosite single crystal (black lines), compared to calculated linear dichroism corresponding to the isotropic simulated spectra of Fig. 3, scaled in intensity to fit the experimental spectra by factors ranging from 1/2 at 290 K down to 1/3 for the 60 and 40 K data. Possible reasons for this mismatch are discussed in the text.

applying an external field, the x-ray magnetic circular dichroism (XMCD) and x-ray magnetic linear dichroism (XMLD) sum rules can be applied to obtain the expectation values of the orbital and spin magnetic moments⁴⁶ and the magneto-crystalline anisotropy⁴⁷, respectively. In the case of the jarosites it is not possible to align the magnetic moments parallel with an external field due to the large energy scale of the antiferromagnetic interactions ($\Theta_{\text{CW}} \approx -800$ K). The only way to obtain accurate information about the relevant expectation values is by explicitly calculating the transition probabilities

$$f(E_{\hbar\omega}, \mathbf{q}) = \frac{|\langle \psi_g | \hat{\mathbf{r}}_{\mathbf{q}} | \psi_e \rangle|^2}{E_{\hbar\omega} - \Delta E - i\Gamma/2}, \quad (2)$$

between the ground state ψ_g and excited state ψ_e , where $\hat{\mathbf{r}}_{\mathbf{q}}$ is the electric-dipole operator for x-ray absorption with polarization \mathbf{q} , and Γ is the lifetime broadening, which is treated as a fitting parameter. The experimental x-ray spectra were used to obtain good approximations for ψ_g and ψ_e with the aid of the Hilbert++ code^{48,49}. The Hilbert space for $|\psi_g\rangle$ is spanned by deter-

minants of the lowest atomic LSJ configurations of $|3d^n\rangle$ and $|3d^{n+1}\underline{L}\rangle$, where \underline{L} denotes a ligand hole. Likewise, the Hilbert space for $|\psi_e\rangle$ is spanned by determinants of the atomic configurations of $|2p^53d^{n+1}\rangle$ and $|2p^53d^{n+2}\underline{L}\rangle$. The Slater integrals for the $3d$ - $3d$ and $2p$ - $3d$ interactions and the $2p$ and $3d$ spin-orbit parameters for the ground state and final state configurations calculated using Cowan's atomic multiplet program⁵⁰ are tabulated in Ref. 51. Approximate solutions for the eigenfunctions of the iron ion in the compound are then obtained by diagonalization of the Hamiltonian including charge-transfer and crystal-field terms arising from the surrounding oxygen ligands⁴⁹

$$\begin{aligned} \mathcal{H} = & \mathcal{H}_{\text{atom}} \\ & + \sum_b [t_{\sigma,b} d_{3\tilde{z}^2-r^2}^+ p_{\tilde{z}}^- + t_{\pi,b} (d_{\tilde{x}\tilde{z}}^+ p_{\tilde{x}}^- + d_{\tilde{y}\tilde{z}}^+ p_{\tilde{y}}^-) + \text{cc}] \\ & + \sum_b [V_{\sigma,b} d_{3\tilde{z}^2-r^2}^+ d_{3\tilde{z}^2-r^2}^- + V_{\pi,b} (d_{\tilde{x}\tilde{z}}^+ d_{\tilde{x}\tilde{z}}^- + d_{\tilde{y}\tilde{z}}^+ d_{\tilde{y}\tilde{z}}^-)] \\ & + \epsilon_p \sum_b (p_{\tilde{x}}^+ p_{\tilde{x}}^- + p_{\tilde{y}}^+ p_{\tilde{y}}^- + p_{\tilde{z}}^+ p_{\tilde{z}}^-). \end{aligned} \quad (3)$$

In this Hamiltonian, the first term represents the atomic Hamiltonian. The second term gives the hybridization between metal $3d$ and oxygen $2p$ orbitals, where $t_{\sigma,b}$ and $t_{\pi,b}$ are the Slater-Koster hopping parameters. For each bond, $t_{\sigma,\pi}$ are rescaled with the bond length using $t_{\sigma,\pi,b} = t_{\sigma,\pi} (R_b/R_{\text{ref}})^\alpha$ where R_b/R_{ref} is the normalized length of Fe-O bond b and the rescaling exponent α is treated as a free parameter. The third term gives the contribution of the electrostatic crystal field which was not used in our case. The last term gives the energy of the $2p$ electrons in the oxygen valence band, where ϵ_p is a simulation parameter related to the charge-transfer gap Δ_{pd} via $\Delta_{pd} = U_{3d^5} - \epsilon_p$. The total dd Coulomb energy for the addition of a sixth electron in the Fe $3d$ shell $U_{3d^5} = \sum_{n=1}^5 U_n$ was found to be 44.3 eV by considering the case $t_{\sigma,\pi} = 0$, where no charge transfer occurs unless the $\text{Fe}^{2+} d^6 \underline{L}$ configuration has a lower energy than the $\text{Fe}^{3+} d^5$ configuration. The sum is made over all bonds b , where the local coordinate frame $\tilde{x}, \tilde{y}, \tilde{z}$ is oriented with the \tilde{z} axis along the bond direction. Numerous comparisons between experiment and theory^{41,52} have shown that this approach works well, provided that the Slater integrals for the Coulomb interactions and spin-orbit couplings are reduced to 70-80% of the atomic Hartree-Fock values³². Here, all pd and dd Slater integrals were reduced to 80% of the atomic values calculated using Cowan's program. The lifetime broadening Γ was set to 0.24 eV for the L_3 edge and 0.3 eV for the L_2 edge. These are values as commonly used in the calculation of x-ray absorption spectra, with L_2 always slightly larger than L_3 . This spectral broadening is due to the coupling with the (infinite dimensional) electromagnetic field.

The crystal field at the iron atom in potassium iron jarosite is approximately of D_{4h} symmetry, with an elongation of the bond length R_{ap} along the fourfold symmetry axis (c') of $\sim 4\%$ with respect to the bond length R_{eq}

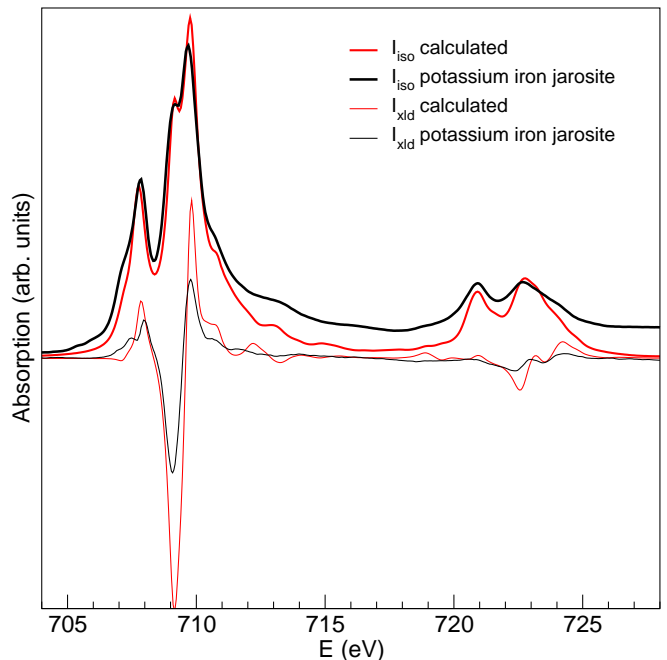


FIG. 5: (color online) Experimentally obtained isotropic spectrum and x-ray linear dichroism (black lines) at 290 K and simulated XAS and XLD (red lines) Boltzmann averaged at 110 K. The simulated XLD is roughly twice as large as the experimentally obtained XLD which is ascribed to misalignments and sample imperfections.

in the equatorial plane¹⁶. The local symmetry axes of the FeO_6 octahedra are canted with respect to the crystallographic c axis by an angle of 19.2° , resulting in buckled FeO_4 planes forming the kagome layers.

Our simulation shows that the spectra are mainly determined by the strongly anisotropic σ -type hybridization between the iron $3d$ and oxygen $2p$ orbitals. The electrostatic part of the crystal field could be left out. An excellent agreement between experimental and simulated spectra for I_{iso} was obtained using only four simulation parameters, namely the charge-transfer gap Δ_{pd} , t_σ , t_π and the rescaling constant α as defined in $t_{\pi,\sigma}^{\text{ap}} = t_{\pi,\sigma}^{\text{eq}} (R_{\text{ap}}/R_{\text{eq}})^\alpha$ (with $R_{\text{ap}}/R_{\text{eq}} = 1.035$ in the present case). When the simulated isotropic spectra is in good agreement with the experimental spectrum, at the same time optimal agreement is obtained for the linear dichroism spectrum. At high temperature there is good agreement between the experimental and simulated linear dichroism spectrum, apart from an intensity scaling factor of 0.5 as shown in Figs. 4 and 5. With the fitting parameters for the simulation there is no way to reduce the overall intensity of the linear dichroism without radically modifying the isotropic spectrum. For up to 6% the intensity difference might be attributable to the local canting of the FeO_6 octahedra with respect to the ab -plane. In addition to this there is an inevitable loss of intensity in the XLD due to misalignments and sample

imperfection. Reduced energy resolution can also lead to a reduction of the measured XLD. We conclude that the difference in overall intensity of the experimental and calculated I_{xld} is not due to fundamental shortcomings of the simulation. For reasons we do not quite understand, the agreement between theory and experiment as evident in the temperature dependence of the isotropic spectra is not reflected in the XLD spectra at lower temperatures. We expect that these spectra, which are averages over increasingly fewer states, are also increasingly sensitive to shortcomings in our simulations. Since the XLD is an orientational energy dependence it will be particularly sensitive for any shortcomings in our theoretical model. It should be emphasised that, despite the differences between the experimental and simulated spectra at low temperatures, the overall agreement between experiment and simulation as obtained here is remarkably good. The parameters as found in our simulation are given in Table I (left-hand side). The large error bar on the value for the charge-transfer gap of 3.5(2.0) eV reflects the adjustment of the total dd Coulomb energy on the iron site $U_{dd}^{\text{tot}} = U_{3d^5}(\langle n \rangle - 5)$ to any increase in ϵ_p by transfer of additional electron weight into the Fe ion. This is reflected in the total error in $\langle n \rangle$ of ~ 0.1 electron.

The values found here for Δ_{pd} and t_σ are similar to those found previously for other Fe^{3+} compounds^{41,53}. Fe_2O_3 , which is in many respects probably the compound most closely related to iron jarosite, was previously found to have a band gap of predominantly charge-transfer character of 2 to 2.5 eV⁵⁴. With an unusually small value for t_π of 0.3(2) eV the lowest energy peak, at 707.8 eV is reproduced in the simulated spectrum. The main peak at 709 eV splits into two peaks with the introduction of a trigonal distortion of the FeO_6 octahedron. A common value for the rescaling exponent α is -3 and the small value found here (-7.5) is an indication that the charge anisotropy is larger than the crystal field anisotropy as defined by the oxygen positions around the iron. The lowest energy low-spin configuration, corresponding to the ${}^4F_{11/2}$ free-ion configuration, lies 1.2 eV higher in energy and is not found to mix into the ground state.

The single-ion anisotropy D_z is obtained as the difference between the energy of the $S_z = \pm 1/2$ and $\pm 5/2$ levels and was found to be of the easy-plane type. This easy-plane anisotropy was also confirmed by rotating the spin quantisation axis to various directions with respect to the crystal field. In this way we have also attempted to obtain estimates for E_{xy} . The D_{4h} symmetry of the ligand field was lowered to C_2 , to model the actual positions of the oxygen anions around the cation as obtained from crystallographical data^{11,16}. However, this had no effect on the agreement of the calculated spectra with experiment, nor on the values as listed in Table I and hence we consider an estimate for E_{xy} from our data presently not attainable. It should in this respect be noted that the value for D_z was found to change with the splitting of the main peak at 709 meV, introduced by the trigo-

TABLE I: Results from simulations of the Fe $L_{2,3}$ spectra of potassium iron jarosite, with ligand-field multiplet calculations using the Hilbert++ code⁴⁹. The left-hand side lists all the parameters varied in the simulation. The right-hand side gives the resulting expectation values for the calculated ground-state properties of the Fe ion. The errors indicate the range over which the parameters could be varied whilst maintaining a good agreement with the data. The right-hand side of the lowest panel gives the effective moment in the case of full orbital polarization, while the calculation indicates only a very weak orbital polarization of the added electron weight due to charge transfer. We use the convention that upper case J, L, S represent operators and lower case j, l, s represent the quantum numbers given by $\langle L^2 \rangle = l(l+1)$ etc.

Simulation pars.		Expectation values	
Δ_{pd} (eV)	3.5(2.0)	D_z (meV)	0.5(1)
t_π (eV)	0.3(2)	$\sqrt{\langle S^2 \rangle} = \langle S \rangle$	2.77(6)
t_σ (eV)	3.0(2)	$\rightsquigarrow s = 2.32$	
α	-7.5	$\sqrt{\langle L^2 \rangle} (\neq \langle L \rangle)$	1.5(2)
		$\rightsquigarrow l = 1.11$	
		but $\langle L \rangle$	0.008(3)
		$\langle n \rangle$	5.39(7)
From $t_{\sigma,\pi}$ and α : (in eV)		Using $\langle L \rangle$ we obtain	
t_σ^{eq}	3.0(2)	$j = s(+l \approx 0)$	2.3
t_σ^{ap}	2.3(2)	g_J	~ 2
t_π^{eq}	< 0.5	μ_{eff}	5.5(1) μ_B
t_π^{ap}	< 0.4	and using $\sqrt{\langle L^2 \rangle}$	
		$j = s + l$	3.4
		g_J	1.67(5)
		μ_{eff}	6.51(7) μ_B

nal distortion of the ligand field. Its value is therefore tightly constrained by the experimentally obtained spectra and the 3d spin-orbit coupling strength. As is detailed in Table I our calculation indicates that for the isolated FeO_6 cluster most of the orbital angular momentum of the Fe $3d^6$ configuration is quenched by the ligand field, with $\langle L \rangle = 0.008(3)$ while the maximum orbital angular momentum which could in this case be realised as $\sqrt{\langle L^2 \rangle} = 1.5(2)$.

Due to the lower resolution at which the hydronium jarosite powder spectrum has been measured we cannot provide detailed information about this system, apart from a confirmation that the general crystal field is similar, with comparable charge-transfer and (σ -type) hybridization. This is a good indication that the magnetic exchange in hydronium jarosite is likely to be comparable with that of the other jarosites, i.e. with a Weiss temperature around -800 K. There is also a slight evidence that the D_{4h} ligand-field anisotropy and hence the zero-field splitting, is slightly smaller in hydronium jarosite.

V. DISCUSSION

The expectation values of the ground-state obtained in our calculation are given in Table I (right-hand side). The corresponding free-ion configuration is $d^5 {}^6S_{5/2}$ with

an additional electron weight of ~ 0.4 electron that is mainly in the $d(x^2 - y^2)$ orbital. The spin-orbit coupling is much smaller than the t_σ hopping bandwidth and hence the fraction of the electron transferred to the ligands to the Fe cation is only very weakly orbitally polarized, with $\langle L \rangle = 0.008(3)$. This value is much smaller than the orbital angular momentum as inferred from high-temperature magnetic susceptibility data, which indicate that $\mu_{\text{eff}} \approx 6.5 \mu_B$ ¹³. As shown in the lower panel of Table I (right-hand side) the experimental value is remarkably close to the case where the added electron weight due to charge transfer is fully orbitally polarized. It is not expected that in this 3d transition metal compound l and $j(=s+l)$ are good quantum numbers, but there is sufficient evidence that $\langle L \rangle$ is significantly larger than zero⁵⁷. It might be that comparison of experimental results with the Mott-Hubbard-like calculations as described here, but with larger clusters containing two or three iron centres, can provide an explanation for the discrepancy between the calculated and experimentally obtained orbital angular momentum.

The energy of $D_z = 0.5(1)$ meV for the easy-plane single-ion anisotropy is in good agreement with the values found by Matan *et al.*²³ (0.428(5) meV) and Coomer *et al.*²⁴ (0.47(2) meV) in their spin-wave analysis of inelastic neutron data from single crystals and powders, respectively. This implies that an easy-plane anisotropy, along with further neighbor interactions, can explain the magnetic ground state in potassium iron jarosite. Further evidence might be found in the temperature dependence of the experimental and calculated isotropic spectra. The spectra in Fig. 5 were taken at room temperature and were compared with the calculated spectra for the Boltzmann averaged configurations of the FeO_6 cluster. As is visible in Fig. 3 the relative intensity of the two maxima in the split peak at 709 eV changes in a similar way as was observed in the experimentally obtained spectra. At the same time an (unexplained) increase in intensity of the peak at 707.8 eV peak relative to the intensity of the main peak is observed in the experimental spectra. We have not been able to account for the latter (which was also observed in powder samples) either by changing the temperature of the calculated spectra or with simulations with new values for δ_{pd} , $t_{\sigma,\pi}$ and α . The change in the main peak at 709 eV where we obtain good agreement, is a direct result of the spin-orbit coupling which causes a change in the shape of the electronic 3d orbital as the spins align in the easy plane. The change in shape of the main peak in the experimental spectra was found to occur between 200 K and 40 K. This is in agreement with previous experiments which have shown that in potassium jarosite the alignment of spin with the kagome planes sets in above 120 K, i.e. well above the transition temperature of 65 K¹³. The temperature dependence in the experimentally obtained spectra was simulated with the temperature (given in units of eV) used for the Boltzmann averaging over excited states as the only extra parameter, as shown in Fig. 3. All other simulation parameters

were fixed to the values given in Table I. In this way we found the expectation values for $\langle |S_z| \rangle$ at each corresponding experimental temperature, arising from the thermal population of the m_S levels in the $S = 5/2$ multiplet. The temperature dependence of $\langle |S_z| \rangle$ (Fig. 6) shows a gradual alignment of the spins into the kagome planes, at temperatures well above the transition temperature to a long-range ordered state at 64 K. This is in rough agreement with the degree of co-planarity of the spins as measured using neutron diffraction on large single crystals¹³.

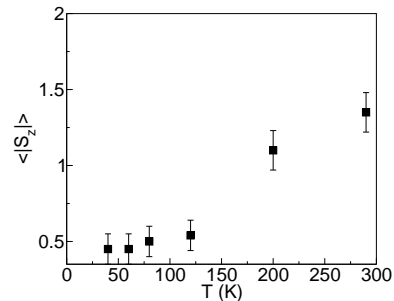


FIG. 6: The expectation values $\langle |S_z| \rangle$ of the simulated spectra, against the temperature of the corresponding experimental spectrum. For a single $S = 5/2$ spin the minimum for $\langle |S_z| \rangle$ is 0.5 and the maximum is 1.5 (equal population of all m_S). The gradual increase of $\langle |S_z| \rangle$ with temperature is in agreement with earlier neutron measurements which show a gradual alignment of the spins in the kagome plane at temperatures well above the transition temperature of 64 K¹³.

The simulated linear dichroism spectra at the lowest temperatures, corresponding to states with the spins aligned in the equatorial plane (the states $|s, m_{z'}\rangle = |2.5, \pm 0.5\rangle$) do not agree well the experimentally obtained spectra as shown in Fig. 4, although some features in the temperature dependence of the XLD are reproduced in the simulated spectra. Without success we have attempted to improve on this situation by calculating the spectra for the Fe spin aligned along the x' axis in the equatorial plane (i.e. $|s, m_{x'}\rangle = |2.5, \pm 2.5\rangle$) which resembles more closely the actual ground state in potassium jarosite. The remaining disagreement could be in part due to experimental error due to sample imperfections and in part due to intrinsic limitations of the simulation.

It might be naive to expect that the ligand-field can be inferred from crystallographic data alone. A first indication for this is the large value of α obtained here, which amplifies the effect of the crystallographic trigonal distortion of the FeO_6 octahedra. Even when the symmetry of the ligand field is further lowered from D_{4h} and accurately modelled on the locations of the oxygen anions around the central iron⁵⁸ neither the calculated spectra nor the simulation parameters change significantly. Therefore, it is possible that the actual charge anisotropy follows more closely the xy plane rather than the equatorial oxygens. The easy-plane magnetic anisotropy of

$D_z = 0.5(1)$ meV found here (see Table I), is then not directly related to the trigonal distortion of the crystal field but more to the 2D character of the electronic hopping and magnetic exchange, the virtual absence of π symmetry hybridization and the Fe $3d$ spin-orbit coupling strength. In that case the easy plane could coincide more closely with the kagome plane. The strong 2D character can in part be attributed to the sulphate groups separating the kagome layers. The formal valence of the sulphur cation is $6+$ with the [Ne] electronic configuration, but this large electrostatic charge is likely to attract electron weight from the neighboring oxygen anions, including the oxygens apical in the FeO_6 octahedra. The hopping parameters $t_{\sigma,\pi}$ from these oxygens to the iron is then further reduced with respect to the in-plane hopping. This translates into the large value for α found here.

The large difference between t_π and t_σ obtained here might be explained by the triangular arrangement of the neighboring FeO_6 octahedra. This dramatically lowers the symmetry of the local crystal field around the oxygen ligands and reduces the π symmetry overlap of the oxygen $2p$ orbitals with the iron $3d$ orbitals. Due to the low symmetry of the oxygen sites the symmetry of the crystal field can be expected to be even lower than indicated by the crystallographic positions of the anions. It is in this sense fortunate that the situation can be modelled so well by the anisotropic π and σ hopping whose ratio is equal for all six anions around the central iron. It could well be that the last differences between the calculated and experimental spectra, such as the low temperature x-ray linear dichroism, some of the difference in the overall intensity of the XLD and the missing bump just above 707 eV, could be brought into agreement by allowing a different balance between t_π and t_σ for the apical oxygen anions.

In neutron spin-wave studies^{21,23,24} the crystal field anisotropy scenario has been compared in close detail with that of DMI. In the present context such a discussion would require calculations on clusters containing two to three iron centres. Such a calculation would certainly be of great interest, but has not been attempted here, since the present approach already accounts so well for the experimental observations on potassium iron jarosite. The present result closes the last missing causal link, from electronic structure to collective magnetic ground state.

VI. CONCLUSION

The isotropic and x-ray linear dichroism spectra taken at the Fe $L_{2,3}$ edges in potassium iron jarosite indicate that an orbital angular momentum and zero-field splitting arise due to an anisotropic, mainly σ type, charge-

transfer from the oxygen ligands. This closely resembles the scenario described in Ref. 36 to explain the zero-field splitting as observed in Mn^{2+} salts.

The zero-field splitting parameter of the crystal field anisotropy of 0.5 meV as found here is in excellent agreement with the value found from fits of a crystal field anisotropy model to the spin-wave spectrum in the ground state of potassium iron jarosite^{21,23,24}. However, the most important result obtained here is that we have verified in detail the mechanism by which this magnetic anisotropy arises, closing the remaining missing link between the electronic structure in potassium iron jarosite and its magnetic ground state. We have even measured a change in the electronic structure as the magnetic ground state becomes co-planar, with the change of the isotropic line shape over the relevant temperature range. Of course, an additional in-plane anisotropy E_{xy} and further-neighbor interactions are needed to fully explain the magnetic ground state. The present study does not provide more hints to the origin of the glassy state in hydronium jarosite. We found that the energy scale of the magnetic interaction in hydronium jarosite should be comparable to that in potassium jarosite. It might be that temperature dependent spectra taken at a higher resolution could provide new insights in the case of hydronium jarosite. A final interesting observation is that the maximum orbital angular momentum polarization possible for the amount of charge-transfer observed here, but not the actually calculated orbital polarization, is in rough agreement with the experimentally observed effective moment.

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- ⁵⁶ It has been shown³⁰ that in 4th, 5th and 6th order perturbation theory of the crystal field and spin-spin interactions a small orbital angular momentum and single-ion anisotropy arises.
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- ⁵⁸ Though all Fe-O_{eq} bond lengths are the same, the equatorial plane is rectangular rather than square, with the short side pointing into the kagome triangles. Furthermore, the Fe-O_{ap} axis is canted by 4° away from the local c' axis.